



# THE MUSCLES TREASURY SURVEY. II. INTRINSIC $\text{Ly}\alpha$ AND EXTREME ULTRAVIOLET SPECTRA OF K AND M DWARFS WITH EXOPLANETS\*

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## ABSTRACT

The ultraviolet (UV) spectral energy distributions (SEDs) of low-mass (K- and M-type) stars play a critical role in the heating and chemistry of exoplanet atmospheres, but are not observationally well-constrained. Direct observations of the intrinsic flux of the  $\text{Ly}\alpha$  line (the dominant source of UV photons from low-mass stars) are challenging, as interstellar H I absorbs the entire line core for even the closest stars. To address the existing gap in empirical constraints on the UV flux of K and M dwarfs, the MUSCLES *Hubble Space Telescope* Treasury Survey has obtained UV observations of 11 nearby M and K dwarfs hosting exoplanets. This paper presents the  $\text{Ly}\alpha$  and extreme-UV spectral reconstructions for the MUSCLES targets. Most targets are optically inactive, but all exhibit significant UV activity. We use a Markov Chain Monte Carlo technique to correct the observed  $\text{Ly}\alpha$  profiles for interstellar absorption, and we employ empirical relations to compute the extreme-UV SED from the intrinsic  $\text{Ly}\alpha$  flux in  $\sim 100$  Å bins from 100–1170 Å. The reconstructed  $\text{Ly}\alpha$  profiles have  $300 \text{ km s}^{-1}$  broad cores, while  $>1\%$  of the total intrinsic  $\text{Ly}\alpha$  flux is measured in extended wings between 300 and  $1200 \text{ km s}^{-1}$ . The  $\text{Ly}\alpha$  surface flux positively correlates with the Mg II surface flux and negatively correlates with the stellar rotation period. Stars with larger  $\text{Ly}\alpha$  surface flux also tend to have larger surface flux in ions formed at higher temperatures, but these correlations remain statistically insignificant in our sample of 11 stars. We also present H I column density measurements for 10 new sightlines through the local interstellar medium.

*Key words:* ISM: clouds – stars: low-mass

*Supporting material:* figure set

## 1. INTRODUCTION

Ultraviolet (UV) photons control the upper atmospheric heating and chemistry of exoplanets. For exoplanets around low-mass stars (spectral types K and M), characterization of the UV environment has only recently begun (e.g., Walkowicz et al. 2008; France et al. 2013; Shkolnik et al. 2014). Most prior observations focused on flare stars (Landsman & Simon 1993; Ayres et al. 1995; Hawley et al. 2003; Osten et al. 2005), and atmospheric models of cool stars have mostly focused on photospheric, Ca II H & K, and H $\alpha$  emission (Husser et al. 2013, but also see Fuhrmeister et al. 2005). UV observations of several optically inactive M dwarfs have shown that their chromospheres and transition regions are active and contribute significant flux to the UV spectral energy distributions (SEDs; France et al. 2013). The very few previously available UV spectra of M dwarfs have limited

our ability to accurately model the atmospheres and potential habitability of rocky planets around low-mass stars.

M dwarfs are the most ubiquitous stellar type in the universe, outnumbering G stars 12–1 within 10 pc (The RECONS Team<sup>11</sup>). Exoplanets in M dwarf habitable zones are easier to detect by both the transit and radial velocity methods, making targeted searches for habitable exoplanets around M dwarfs potentially more fruitful than for any other stellar type. Dressing & Charbonneau (2015) estimated that there are 0.18–0.27 Earth-sized planets ( $1\text{--}1.5 R_{\oplus}$ ) and 0.11–0.25 super-Earths ( $1.5\text{--}2 R_{\oplus}$ ) per M dwarf habitable zone.

Some of the issues thought to threaten the habitability of an exoplanet orbiting a low-mass star are currently believed to be surmountable (e.g., Tarter et al. 2007). Joshi (2003) showed that close-in, tidally locked exoplanets can maintain a habitable climate as long as significant surface pressure exists to support atmospheric heat transport to the nightside. The high-energy flux with which M dwarfs irradiate their exoplanets may be similar to the young Sun (Lammer et al. 2011). Joshi (2003) noted that numerous starspots can cause the stellar luminosity

\* Based on observations made with the NASA/ESA *Hubble Space Telescope*, obtained from the data archive at the Space Telescope Science Institute. STScI is operated by the Association of Universities for Research in Astronomy, Inc. under NASA contract NAS 5-26555.

<sup>11</sup> <http://www.recons.org>

to vary with rotation by approximately 40%, but the same change is experienced by a planet with the eccentricity of Mars, and life may have been possible on Mars (Noffke 2015). Unresolved issues include the effects of particle fluxes (enhanced during flares and coronal mass ejections) on the magnetospheres and atmospheres (e.g., Lammer et al. 2007; Segura et al. 2010; Cohen et al. 2014), UV-driven water loss (Luger et al. 2015), and UV-driven photochemistry in terrestrial and gaseous exoplanet atmospheres (Grenfell et al. 2014; Miguel et al. 2014; Rugheimer et al. 2015).

Observations of potential biosignatures (e.g., O<sub>2</sub>, O<sub>3</sub>, and CH<sub>4</sub>) and habitability indicators (e.g., H<sub>2</sub>O and CO<sub>2</sub>) in these planets' atmospheres will likely be possible in the next two decades (e.g., 2015 AURA High-Definition Space Telescope report<sup>12</sup>). However, UV-driven photochemistry can mimic certain biosignatures, complicating the analysis (Domagal-Goldman et al. 2014; Tian et al. 2014; Harman et al. 2015). Ly $\alpha$  and other bright chromospheric and transition region resonance lines in the far-UV (912–1700 Å) photodissociate CO<sub>2</sub> and H<sub>2</sub>O, resulting in a buildup of abiotic O<sub>2</sub>. O<sub>2</sub> dissociation by 1200–2400 Å photons leads to a three-body reaction involving O and O<sub>2</sub> that accumulates abiotic O<sub>3</sub> in the exoplanet's atmosphere. Because these reactions depend on the relative balance of far-UV to near-UV (1700–3200 Å) flux, characterizing the full stellar UV spectrum is critical for atmospheric modeling to interpret the origin of biosignatures.

Low-mass stars have weak near-UV and far-UV continua punctuated by emission lines from the chromosphere and transition region, including H I Lyman alpha (Ly $\alpha$ ;  $\lambda$ 1216 Å). Ly $\alpha$  comprises  $\sim$ 37%–75% of the total 1150–3100 Å flux from most M dwarfs (France et al. 2013), but is significantly attenuated by optically thick H I absorption from the intervening interstellar medium (ISM), even for the closest stars ( $d < 22$  pc in this study). The intrinsic Ly $\alpha$  emission profile must be reconstructed from the observed profile (e.g., Wood et al. 2005; France et al. 2013; Linsky et al. 2013) to properly characterize the far-UV spectrum.

The Ly $\alpha$  flux is also one proxy used to estimate the extreme-UV (100–912 Å) flux, which heats exoplanetary upper atmospheres, driving exoplanet mass-loss (Lammer et al. 2003, 2007, 2009). The extreme-UV is heavily attenuated for all stars, except the Sun, by neutral hydrogen in the ISM from 912 Å to approximately 400 Å, where ground-state neutral hydrogen's photoionization cross-section becomes an order of magnitude smaller. No astronomical instrument currently exists to observe the 100–400 Å spectral window.

An unknown intrinsic profile shape and complicated ISM structure along the line of sight are challenges when reconstructing the intrinsic Ly $\alpha$  profiles of low-mass stars. The Sun's intrinsic Ly $\alpha$  profile is hundreds of km s<sup>-1</sup> broad with narrow self-absorption at the peak (e.g., Fontenla et al. 1988; Woods et al. 1995; Curdt et al. 2008), and its Mg II profile has a similar shape (Schmit et al. 2015). For K dwarfs, Wood et al. (2005) assumed a Ly $\alpha$  profile similar to the ISM absorption-corrected Mg II line, which exhibits self-reversal. M-dwarf temperature-pressure profiles may be different enough from the Sun's that self-reversal may not be present. Sunspots, which have similar effective temperatures to M dwarfs, do not show Ly $\alpha$  self-reversal (Fontenla et al. 1988; Curdt et al. 2001; Tian et al. 2009), and M-dwarf Mg II profiles do not show self-

reversal (Wood et al. 2005). Also, Ly $\alpha$ 's linewidth is expected to be much narrower than the Sun's due to the positive correlation between stellar luminosity and linewidths of optically thick chromospheric lines (Wilson & Vainu Bappu 1957; Ayres 1979; Cassatella et al. 2001).

The *Hubble Space Telescope* (HST) MUSCLES<sup>13</sup> Treasury Survey has directly measured the 1150–3200 Å spectra of 11 nearby ( $d < 22$  pc) K and M dwarf exoplanet hosts in order to provide an empirical basis for stellar-irradiance estimates. Paper I (France et al. 2016) provides an overview of the science goals and observing strategy of the MUSCLES survey as well as a summary of initial results. In this Paper II of a series from the MUSCLES team, we present the reconstructions of the intrinsic Ly $\alpha$  stellar emission lines and the extreme-UV irradiance. Paper III (Lloyd et al. 2016) discusses the data reduction steps followed to merge spectra from the HST observations and other sources into panchromatic SEDs and examines the photodissociation of several common molecules when directly exposed to these SEDs. In this work, Sections 2 and 3 describe the Ly $\alpha$  observations, reductions, and the methodology we use to reconstruct the intrinsic Ly $\alpha$  profiles for the stars in the survey. Section 4 presents the Ly $\alpha$  and extreme-UV reconstructions that have been incorporated into panchromatic spectra of our targets (available as High Level Science Products on MAST<sup>14</sup>). We compare Ly $\alpha$  with other chromospheric emission lines and discuss the profile shape in Section 5 and summarize our results in Section 6.

## 2. LY $\alpha$ OBSERVATIONS AND REDUCTIONS

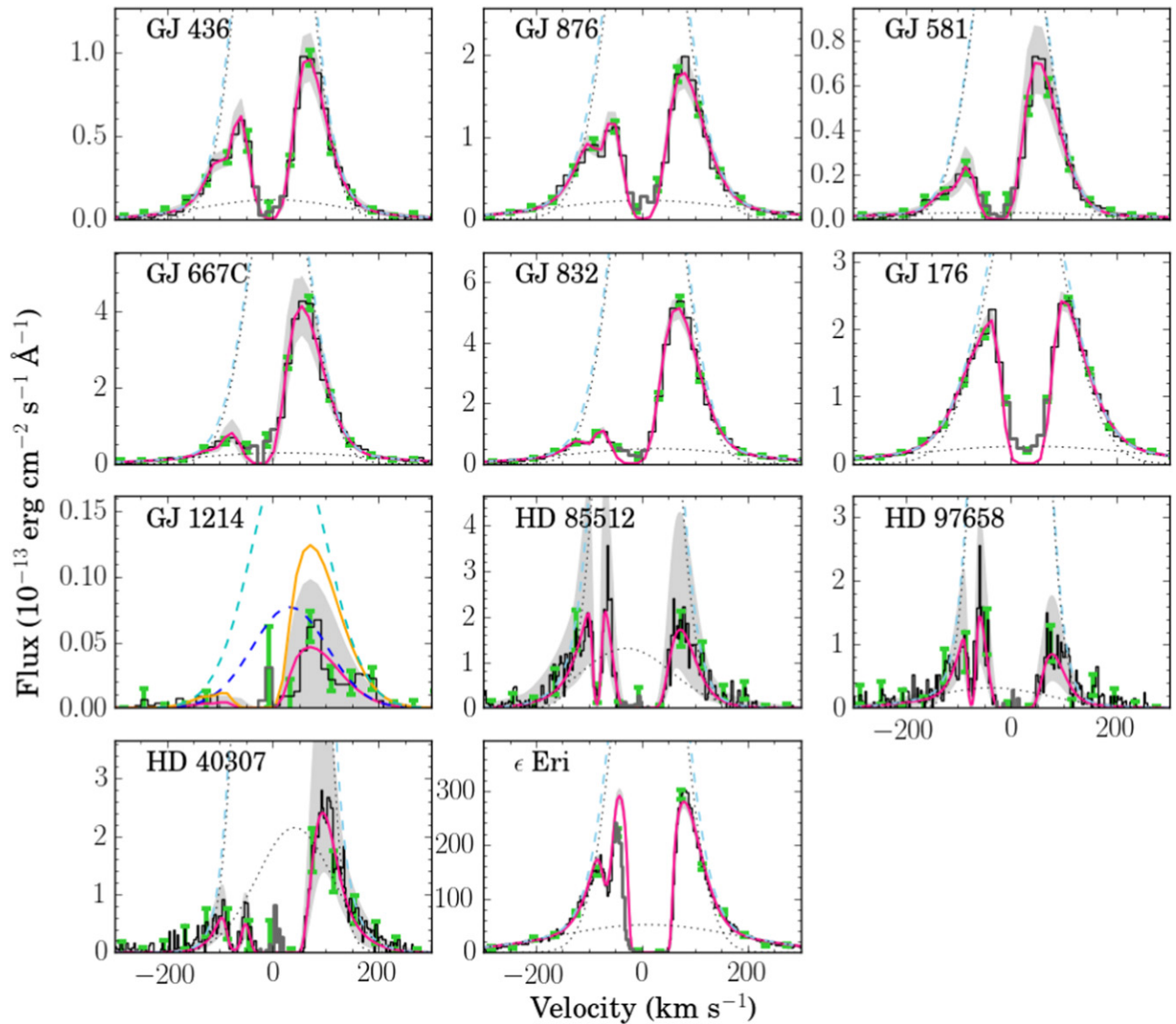
Observations were obtained between 2014 October and 2015 August with the Space Telescope Imaging Spectrograph (STIS) and Cosmic Origins Spectrograph (COS) on HST (see France et al. 2016 for an overview of the observing strategy). To observe Ly $\alpha$ , we used the STIS G140M grating ( $R \sim 11,400$ ) with the  $52'' \times 0''.1$  slit for the seven M dwarfs and the E140M grating ( $R \sim 48,500$ ) with the  $0''.2 \times 0''.06$  slit for the four K dwarfs. The small slit sizes were chosen to minimize geocoronal Ly $\alpha$  airglow contamination. The absolute velocity accuracy near Ly $\alpha$  for G140M is 2.5–6.2 km s<sup>-1</sup> and for E140M is 0.7–1.6 km s<sup>-1</sup>.

Although STIS's narrow slit minimizes airglow photons entering the spectrograph, geocoronal Ly $\alpha$  can still dominate the stellar Ly $\alpha$  signal. The H I column densities of the local ISM are large enough ( $N(\text{H I}) \gg 14$ ) to completely attenuate the Ly $\alpha$  line cores for the stars in our sample ( $\log_{10} F(\text{Ly}\alpha)$  is about  $-14$  to  $-10$  erg cm<sup>-2</sup> s<sup>-1</sup>). To further minimize geocoronal Ly $\alpha$  contamination of the observed Ly $\alpha$  profiles, we scheduled each target's observations during the time of year when Earth's heliocentric velocity would allow the geocoronal Ly $\alpha$  airglow to coincide with the attenuated stellar line core. Airglow was subtracted using the detector regions outside the spectral trace. Residual airglow is resolved in the E140M observations, allowing for a second airglow subtraction from the one-dimensional spectra. The G140M observations do not have sufficient resolution for the second iteration of airglow subtraction, so we mask the attenuated line core from the fits (see Section 3), because the line cores should be consistent with zero flux for all of our targets. To avoid biasing the Ly $\alpha$  line

<sup>12</sup> <http://www.hdstvision.org/release-text>

<sup>13</sup> Measurements of the Ultraviolet Spectral Characteristics of Low-mass Exoplanetary Systems.

<sup>14</sup> <https://archive.stsci.edu/prepds/muscles/>



**Figure 1.** MCMC solutions for the seven M dwarfs and four K dwarfs. The black histograms with green error bars show STIS G140M or E140M data, the pink lines are the best-fit attenuated profiles determined by MCMC, the gray-shaded regions show the uncertainty of the fits, and the dashed blue lines are the reconstructed intrinsic  $\text{Ly}\alpha$  profiles. The dotted lines show the narrow and broad components of the reconstructed profiles. The sections of data masked from the fits are shown with the dark gray histogram. In GJ 1214’s panel, the solid pink line, gray-shaded region, and dashed, dark blue line represent the original MCMC fit, and the solid orange and dashed light blue lines represent the scaled solution (see Section 4.7).

profile fits to the low-signal wings, we ensured the error bars were greater than or equal to the rms scatter in the wings over the velocity range  $-300$  to  $-400$   $\text{km s}^{-1}$ .

Because COS’s aperture is large ( $2''5$ ), the COS G130M observations are too contaminated by airglow to access the  $\text{Ly}\alpha$  line core ( $\pm 300$   $\text{km s}^{-1}$ ), but they do detect the broad wings of  $\text{Ly}\alpha$  (see Section 5.1). We refer the reader to Loyd et al. (2016) for details on the reduction of the STIS and COS observations and the construction of the panchromatic SEDs.

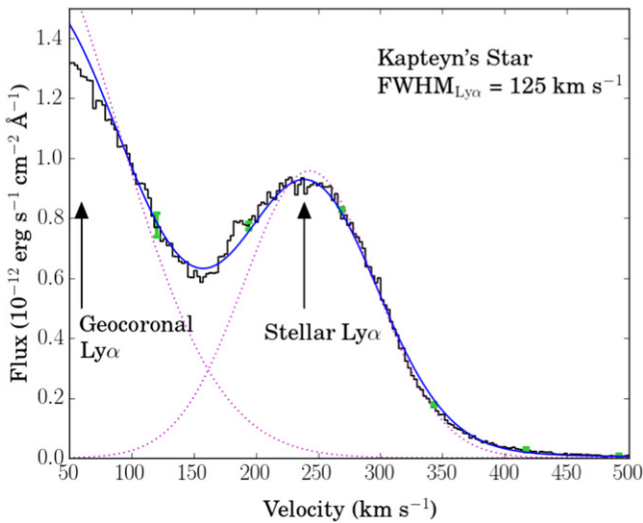
### 3. METHODOLOGY

#### 3.1. The Model

To reconstruct the intrinsic  $\text{Ly}\alpha$  profiles from the observed profiles (presented fully in Section 4), we fit a nine-parameter model describing the intrinsic stellar  $\text{Ly}\alpha$  profile (six parameters) and the interstellar H I and D I absorption (three parameters) via a Markov Chain Monte Carlo (MCMC)

technique (see Section 3.2). Figure 1 shows the observed  $\text{Ly}\alpha$  profiles (black histograms with green error bars). In  $\epsilon$  Eri’s panel, the H I and D I absorption are seen around  $0$   $\text{km s}^{-1}$  and  $-75$   $\text{km s}^{-1}$ , respectively. The fits shown in Figure 1 will be discussed in detail in Section 4.

We assume for the intrinsic  $\text{Ly}\alpha$  profile a dominant narrow Gaussian and weaker broad Gaussian, each described by an amplitude  $A$  ( $\text{erg cm}^{-2} \text{s}^{-1} \text{\AA}^{-1}$ ), FWHM ( $\text{km s}^{-1}$ ), and heliocentric velocity centroid  $V$  ( $\text{km s}^{-1}$ ). The  $\text{Ly}\alpha$  line has not been modeled from first principles for any M dwarf, and semi-empirical models provide only a rough fit to the data (Fontenla et al. 2016). As discussed in Section 1, no self-reversal of the line core is expected. We searched the *HST* archive for a  $\text{Ly}\alpha$  observation of a high radial velocity M dwarf whose stellar emission is Doppler shifted away from the geocoronal and ISM line centers. We analyzed the *HST*–COS spectrum of Kapteyn’s Star (Guinan et al. 2016; GJ 191; spectral type sdM1.0). Figure 2 shows its Gaussian-shaped  $\text{Ly}\alpha$  profile with



**Figure 2.** COS G130M spectrum (Guinan et al. 2016) of Kapteyn’s star (spectral type sdM1.0) is shown with the black histogram and green error bars. The blue line is the best simultaneous fit to the wing of the geocoronal Ly $\alpha$  and stellar Ly $\alpha$  emission, with magenta dotted lines showing the individual components of the fit before convolution with the wavelength-dependent COS line-spread function. The FWHM of Kapteyn’s Star’s Ly $\alpha$  profile is 125 km s $^{-1}$ .

no self-reversal, redshifted ( $V = +245$  km s $^{-1}$ ) well away from the geocoronal emission and ISM absorption. This supports our choice of a double-Gaussian as the intrinsic Ly $\alpha$  profile shape for the M dwarfs.

We assume for the K dwarfs the same profile shape as for the M dwarfs. Wood et al. (2005) used the absorption-corrected Mg II profiles, which display self-reversal, as a Ly $\alpha$  template. In the absence of a Mg II template, Bourrier et al. (2013) used two Voigt profiles of equal width and damping constants separated in velocity as the intrinsic Ly $\alpha$  profile. The velocity separation of the two emission components can mimic self-reversal of the line core. We discuss the effect of excluding self-reversal on the K-dwarf Ly $\alpha$  reconstructions in Section 3.3.1.

To model the ISM H I and D I absorption, we assume a single Voigt profile parameterized by H I column density  $N(\text{H I})$  (cm $^{-2}$ ), Doppler  $b$  value (km s $^{-1}$ ), velocity centroid  $V_{\text{H I}}$  (km s $^{-1}$ ), and D/H ratio fixed at  $1.5 \times 10^{-5}$  in accord with the mean value found by Wood et al. (2004) for the local ISM. We assume that thermal line broadening determines the Doppler width parameter ( $b_{\text{D I}} = b_{\text{H I}}/\sqrt{2}$ ), D I has the same velocity centroid as H I, and the D I column density scales with the H I column density by the D/H ratio. Local ISM conditions are fairly well known (e.g., Redfield & Linsky 2008), and we expect Doppler  $b$  values around 10 km s $^{-1}$  and  $\log_{10}$  H I column densities around 18 for nearby stars (Wood et al. 2005).

### 3.2. MCMC

We employ the affine-invariant MCMC sampler of Goodman & Weare (2010) via the Python module *emcee* (Foreman-Mackey et al. 2013) to efficiently sample the posterior distribution of our nine parameters (see Table 1 for a list of the parameters). Several of the parameters (e.g., amplitude and column density) are expected to be correlated, and the affine-invariance of *emcee* ensures that the sampling is unaffected by parameter correlations. We assume a Gaussian likelihood and uniform priors for all parameters except the Doppler  $b$

parameter, which has a logarithmic prior probability. We sample the posterior probability distribution 500,000 times for each target with a typical burn-in period of 300 steps and use one- and two-dimensional projections of the posterior probability distributions of the nine parameters (referred to as marginalized and joint distributions, respectively) to analyze the fits. Figure 3 exemplifies the sampled posterior distributions of our high-S/N data (GJ 832; top panel) and that of the low-S/N data (HD 85512; bottom panel). Similar figures for the other nine targets are available in the figure set in the online journal.

Most of the marginalized distributions of the individual parameters are singly peaked, and the joint distributions exhibit the degeneracies between parameters (diagonally aligned contours in Figure 3). The strongest degeneracies are between the narrow component’s amplitude ( $A_{\text{narrow}}$ ) and  $\log_{10} N(\text{H I})$  and between the broad component’s amplitude ( $A_{\text{broad}}$ ) and  $\text{FWHM}_{\text{broad}}$ . Constant intrinsic flux can be maintained by simultaneously increasing (decreasing)  $A_{\text{narrow}}$  and increasing (decreasing)  $\log_{10} N(\text{H I})$ , but the degeneracy between  $A_{\text{broad}}$  and  $\log_{10} N(\text{H I})$  is much weaker. There are also several parameter degeneracies that exhibit the trade-off between having a large or small integrated-flux contrast between the narrow and broad components. Other degeneracies involve the effects of having the velocity centroid of the ISM absorption ( $V_{\text{H I}}$ ) offset from or coincident with the emission peak.

For each of the nine parameters, the best-fit values are taken as the 50th percentile (the median) of the marginalized distributions, and  $1\sigma$  error bars as the 16th and 84th percentiles (shown as dashed vertical lines in Figure 3). The best-fit reconstructed Ly $\alpha$  fluxes are determined from the best-fit amplitude, FWHM, and velocity centroid parameters, and the  $1\sigma$  error bars of the reconstructed Ly $\alpha$  flux are taken by varying these parameters individually between their  $1\sigma$  error bars and keeping all others fixed at their best-fit value. The resulting minimum and maximum Ly $\alpha$  fluxes become the  $1\sigma$  error bars (Table 2).

### 3.3. Systematic Effects

Here we evaluate potential systematic uncertainties not accounted for by the  $1\sigma$  uncertainties determined by the MCMC method. We explore the effects of excluding Ly $\alpha$  self-reversal, assuming a fixed D/H ratio, assuming a single ISM absorption component, and instrumental spectral resolution on the reconstructed Ly $\alpha$  fluxes and the MCMC-derived  $1\sigma$  uncertainties.

The true Ly $\alpha$  flux is not known for any of our target stars. Therefore, we carried out a number of simulations to confirm that our Ly $\alpha$  reconstruction method works and investigated systematic uncertainties potentially not accounted for by the  $1\sigma$  uncertainties determined by the MCMC method. Specifically, we simulated G140M and E140M data assuming different S/N ratios, one to three ISM components, and different D/H ratios. Fitting our model (one ISM component, fixed D/H ratio, no self-reversal) to these data, we find that the reconstructed Ly $\alpha$  flux is always reasonably close to the true Ly $\alpha$  flux and that for all targets except  $\epsilon$  Eri, the uncertainties printed in Table 2 are close to  $1\sigma$ .

We have determined that fits to data with higher instrumental spectral resolution result in  $1\sigma$  uncertainties that do not fully account for systematic uncertainties from the exclusion of self-reversal and additional ISM absorption components (see

**Table 1**  
Effect of Resolution on the MCMC Solution Parameters<sup>a</sup>

	$\epsilon$ Eri		AU Mic		70 Oph A	
	High Res.	Low Res.	High Res.	Low Res.	High Res.	Low Res.
$F(\text{Ly}\alpha) \times 10^{-11}$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$6.1^{+0.2}_{-0.2}$	$5.2^{+1.3}_{-1.0}$	$1.07^{+0.04}_{-0.04}$	$0.9^{+0.2}_{-0.2}$	$3.1^{+0.1}_{-0.1}$	$2.7^{+0.7}_{-0.7}$
<sup>b</sup> $V_{\text{narrow}}$ ( $\text{km s}^{-1}$ )	$13.6^{+0.4}_{-0.4}$	$12.7^{+1.5}_{-1.5}$	$-9.7^{+0.7}_{-0.7}$	$-10.3^{+1.7}_{-1.5}$	$-13.0^{+0.4}_{-0.4}$	$-12.0^{+2.2}_{-2.0}$
<sup>b</sup> $A_{\text{narrow}} \times 10^{-11}$ ( $\text{erg cm}^{-2} \text{s}^{-1} \text{\AA}^{-1}$ )	$9.4^{+0.4}_{-0.4}$	$7.1^{+1.4}_{-1.1}$	$1.2^{+0.06}_{-0.06}$	$0.9^{+0.1}_{-0.1}$	$5.1^{+0.2}_{-0.2}$	$4.2^{+0.8}_{-0.6}$
<sup>b</sup> $\text{FWHM}_{\text{narrow}}$ ( $\text{km s}^{-1}$ )	$129.3^{+1.6}_{-1.6}$	$143.6^{+5.8}_{-5.8}$	$148.6^{+2.7}_{-2.6}$	$180.0^{+11.2}_{-6.4}$	$126.8^{+0.8}_{-0.9}$	$134.0^{+4.3}_{-4.5}$
<sup>b</sup> $V_{\text{broad}}$ ( $\text{km s}^{-1}$ )	$11.7^{+1.8}_{-1.8}$	$11.6^{+9.5}_{-9.3}$	$-8.5^{+1.6}_{-1.5}$	$-4.8^{+5.1}_{-4.8}$	$-4.9^{+1.5}_{-1.6}$	$-5.4^{+6.9}_{-1.6}$
<sup>b</sup> $A_{\text{broad}} \times 10^{-11}$ ( $\text{erg cm}^{-2} \text{s}^{-1} \text{\AA}^{-1}$ )	$0.52^{+0.28}_{-0.26}$	$0.34^{+0.07}_{-0.06}$	$0.19^{+0.01}_{-0.01}$	$0.10^{+0.01}_{-0.03}$	$0.21^{+0.01}_{-0.01}$	$0.15^{+0.03}_{-0.02}$
<sup>b</sup> $\text{FWHM}_{\text{broad}}$ ( $\text{km s}^{-1}$ )	$400.1^{+9.0}_{-8.9}$	$477.9^{+40.6}_{-35.0}$	$387.8^{+9.5}_{-8.7}$	$493.2^{+25.7}_{-23.3}$	$378.4^{+7.3}_{-7.2}$	$431.0^{+29.7}_{-26.2}$
<sup>c</sup> $\log_{10} N(\text{H I})$ ( $\text{cm}^{-2}$ )	$17.93^{+0.02}_{-0.02}$	$17.88^{+0.10}_{-0.11}$	$18.35^{+0.1}_{-0.1}$	$18.32^{+0.03}_{-0.04}$	$18.28^{+0.01}_{-0.01}$	$18.23^{+0.06}_{-0.06}$
<sup>c</sup> $b$ ( $\text{km s}^{-1}$ )	$12.1^{+0.02}_{-0.02}$	$13.8^{+0.4}_{-0.7}$	$14.0^{+0.2}_{-0.2}$	$14.1^{+0.8}_{-0.8}$	$15.5^{+0.1}_{-0.1}$	$15.7^{+0.7}_{-0.1}$
<sup>c</sup> $V_{\text{HI}}$ ( $\text{km s}^{-1}$ )	$16.1^{+0.4}_{-0.4}$	$10.0^{+0.7}_{-0.5}$	$-24.0^{+0.2}_{-0.3}$	$-26.8^{+0.9}_{-1.0}$	$-31.2^{+0.1}_{-0.1}$	$-32.6^{+1.0}_{-1.0}$

#### Notes.

<sup>a</sup> Error bars on the individual parameters are marginalized over all other parameters.

<sup>b</sup>  $\text{Ly}\alpha$  emission parameters.

<sup>c</sup> ISM absorption parameters.

Sections 3.3.1 and 3.3.3 for more details). We find that each can introduce up to a 30% uncertainty. To show the worst case scenario where these uncertainties are present (where appropriate) and at their maximum, we add them here in quadrature with the MCMC  $1\sigma$  uncertainties from Table 2. We add a 30% uncertainty term for excluding self-reversal for all four K dwarfs and another 30% term for HD 85512 and HD 97658 for excluding additional ISM components. Both sightlines are believed to have 2 or 3 ISM components (Redfield & Linsky 2008).  $\epsilon$  Eri's error bars increase dramatically ( $\pm 3\%$  to  $\pm 30\%$ ), and the other three K dwarfs' increase modestly (roughly  $^{+110}_{-45}\%$  to  $^{+115}_{-60}\%$ ). Note that we present the MCMC-determined  $1\sigma$  uncertainties in the reconstructed  $\text{Ly}\alpha$  fluxes in Table 2; these additional systematic terms are not included.

For the G140M observations, we have found some evidence for a  $\sim 15\%$  systematic underestimate of the true intrinsic flux. However, this underestimate is usually less than the  $1\sigma$  error bars on the reconstructed fluxes, so we do not attempt to correct this effect (see Section 3.3.4 for more details).

#### 3.3.1. Excluding Self-reversal

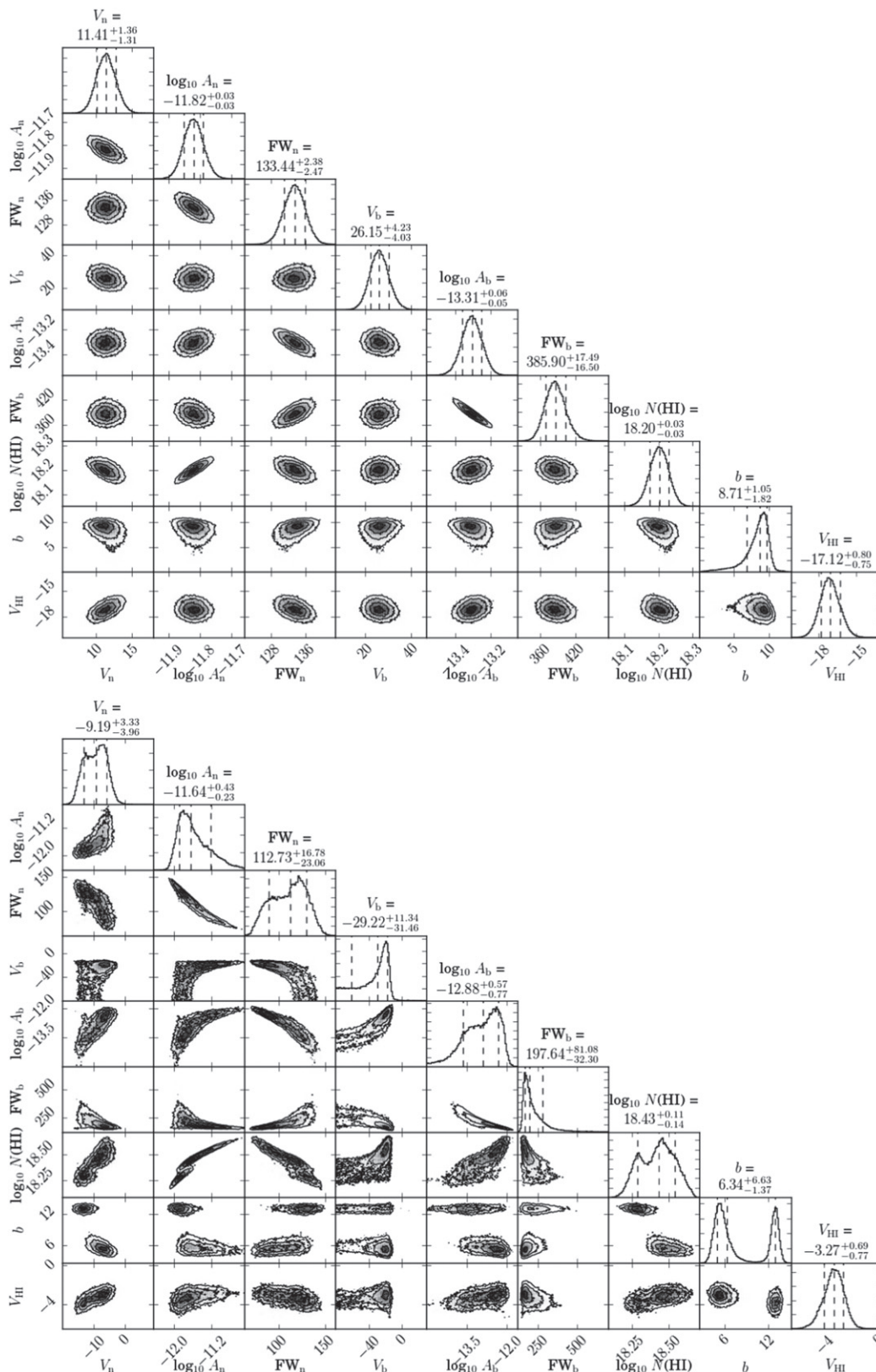
Our model assumes that the K-dwarf intrinsic profiles do not show self-reversal of the  $\text{Ly}\alpha$  line core. While it has not been definitively shown that K stars'  $\text{Ly}\alpha$  profiles are self-reversed, previous authors (Wood et al. 2005; Bourrier et al. 2013) fit self-reversed profiles to observed K-dwarf profiles, so we wish to estimate the effect of excluding self-reversal. We apply our model and methodology to the STIS E140H spectra ( $R \sim 114,000$ ) presented in Wood et al. (2005) for the early K stars  $\epsilon$  Eri, HR 8, and HR 1925. The MCMC reconstructed  $\text{Ly}\alpha$  fluxes for HR 8 and HR 1925 are  $\sim 20\%$  smaller than Wood et al.'s (2005) values. For both stars we find H I column densities

$15\%$ – $20\%$  smaller, thus our reconstructed fluxes are also smaller (see discussion of degeneracies in Section 3.2). For  $\epsilon$  Eri, we find a H I column density 20% larger than Wood et al. (2005) and Dring et al. (1997), thus we find a 27% larger reconstructed  $\text{Ly}\alpha$  flux. Excluding self-reversal appears to introduce up to a 30% difference in the reconstructed  $\text{Ly}\alpha$  fluxes.

#### 3.3.2. Variable D/H Ratio

For the results presented in the next section, we keep the D/H ratio fixed at  $1.5 \times 10^{-5}$  in accord with the average value for the local ISM (Wood et al. 2004), but here we explore the effect of allowing the D/H ratio to be a free parameter with a uniform prior probability. When we allow the D/H ratio to vary (excluding GJ 1214 because its low S/N provides no constraint on the D/H ratio), we find an average D/H ratio in the MUSCLES sample of  $(1.4 \pm 0.6) \times 10^{-5}$ . While this result agrees with Wood et al. (2004), the D/H ratio in these fits was typically not well constrained and had large uncertainties. Except for  $\epsilon$  Eri, the intrinsic  $\text{Ly}\alpha$  fluxes found while floating the D/H ratio all have solutions that are well within the  $1\sigma$  error bars of the intrinsic fluxes found from the fixed D/H ratio fits.  $\epsilon$  Eri's D/H ratio was found to be  $2.2 \times 10^{-5}$  and its reconstructed  $\text{Ly}\alpha$  flux was 13% smaller.

Simulating G140M and E140M observations with different “true” D/H ratios ( $1.3$ – $1.7 \times 10^{-5}$ ) introduces a dispersion in the reconstructed  $\text{Ly}\alpha$  fluxes that is smaller than the dispersion seen from simulating the observations repeatedly. Because the reconstructed  $\text{Ly}\alpha$  flux does not appear to depend sensitively on the assumed D/H ratio, and there is no evidence that the D/H ratio of the local ISM varies significantly from



**Figure 3.** One- and two-dimensional projections of the sampled posterior probability distributions, referred to as marginalized and joint distributions, respectively, of the nine parameters for GJ 832 (top) and HD 8512 (bottom). Contours in the joint distributions are shown at 0.5-, 1-, 1.5-, and 2 $\sigma$ , and the histograms' dashed vertical lines show the 16th, 50th, and 84th percentiles of the samples in each marginalized distribution.

(The complete figure set (11 images) is available in the online journal.)

**Table 2**  
Emission Line Fluxes of MUSCLES Targets

Target Star	$d$ (pc)	$R$ ( $R_{\odot}$ )	$F(\text{Ly}\alpha)^a$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{EUV})^b$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{Mg II})^c$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{Si III})$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{C II})^d$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{Si IV})$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{He II})$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )	$F(\text{C IV})$ ( $\text{erg cm}^{-2} \text{s}^{-1}$ )
$\epsilon$ Eri	3.2	0.74 <sup>t</sup>	$(6.1^{+0.2}_{-0.2}) \times 10^{-11}$	$5.3 \times 10^{-11}$	$(6.28 \pm 0.02) \times 10^{-11}$	$(3.7 \pm 0.2) \times 10^{-13}$	$(4.51 \pm 0.01) \times 10^{-13}$	$(3.12 \pm 0.01) \times 10^{-13}$	$(5.57 \pm 0.03) \times 10^{-13}$	$(7.19 \pm 0.03) \times 10^{-13}$
HD 85512	11.2	0.7 <sup>u</sup>	$(1.2^{+1.8}_{-0.5}) \times 10^{-12}$	$6.8 \times 10^{-13}$	$(6.0 \pm 0.1) \times 10^{-13}$	$(1.7 \pm 0.1) \times 10^{-15}$	$(4.0 \pm 0.1) \times 10^{-15}$	$(1.6 \pm 0.1) \times 10^{-15}$	$(3.0 \pm 0.9) \times 10^{-15}$	$(5.1 \pm 0.8) \times 10^{-15}$
HD 40307	13.0	0.83 <sup>u</sup>	$(2.0^{+2.3}_{-0.9}) \times 10^{-12}$	$1.5 \times 10^{-12}$	$(4.4 \pm 0.1) \times 10^{-13}$	$(1.7 \pm 0.1) \times 10^{-15}$	$(2.8 \pm 0.1) \times 10^{-15}$	$(1.3 \pm 0.1) \times 10^{-15}$	$(2.3 \pm 0.9) \times 10^{-15}$	$(3.0 \pm 0.6) \times 10^{-15}$
HD 97658	21.1	0.72 <sup>v</sup>	$(9.1^{+9.9}_{-4.5}) \times 10^{-13}$	$6.7 \times 10^{-13}$	$(3.08 \pm 0.05) \times 10^{-13}$	$(1.8 \pm 0.1) \times 10^{-15}$	$(2.4 \pm 0.1) \times 10^{-15}$	$(1.3 \pm 0.1) \times 10^{-15}$	$(2.0 \pm 0.3) \times 10^{-15}$	$(2.6 \pm 0.3) \times 10^{-15}$
GJ 832	5.0	0.46 <sup>w</sup>	$(9.5^{+0.6}_{-0.6}) \times 10^{-13}$	$9.6 \times 10^{-13}$	$(2.39 \pm 0.03) \times 10^{-13}$	$(2.5 \pm 0.1) \times 10^{-15}$	$(3.8 \pm 0.1) \times 10^{-15}$	$(3.2 \pm 0.1) \times 10^{-15}$	$(6.8 \pm 0.3) \times 10^{-15}$	$(7.5 \pm 0.3) \times 10^{-15}$
GJ 876	4.7	0.37 <sup>x</sup>	$(3.9^{+0.4}_{-0.4}) \times 10^{-13}$	$3.8 \times 10^{-13}$	$(4.0 \pm 0.1) \times 10^{-14}$	$(8.1 \pm 0.1) \times 10^{-15}$	$(1.06 \pm 0.01) \times 10^{-14}$	$(8.0 \pm 0.1) \times 10^{-15}$	$(5.5 \pm 0.3) \times 10^{-15}$	$(2.27 \pm 0.04) \times 10^{-14}$
GJ 581	6.2	0.3 <sup>y</sup>	$(1.1^{+0.3}_{-0.2}) \times 10^{-13}$	$1.1 \times 10^{-13}$	$(2.4 \pm 0.1) \times 10^{-14}$	$(2.9 \pm 0.4) \times 10^{-16}$	$(4.8 \pm 0.4) \times 10^{-16}$	$(4.1 \pm 0.6) \times 10^{-16}$	$(7.2 \pm 1.9) \times 10^{-16}$	$(1.8 \pm 0.2) \times 10^{-15}$
GJ 176	9.3	0.45 <sup>x</sup>	$(3.9^{+0.2}_{-0.2}) \times 10^{-13}$	$4.0 \times 10^{-13}$	$(2.03 \pm 0.02) \times 10^{-13}$	$(2.1 \pm 0.1) \times 10^{-15}$	$(5.4 \pm 0.1) \times 10^{-15}$	$(2.2 \pm 0.1) \times 10^{-15}$	$(6.7 \pm 0.3) \times 10^{-15}$	$(1.22 \pm 0.03) \times 10^{-14}$
GJ 436	10.1	0.45 <sup>x</sup>	$(2.1^{+0.3}_{-0.3}) \times 10^{-13}$	$2.1 \times 10^{-13}$	$(5.0 \pm 0.1) \times 10^{-14}$	$(5.2 \pm 0.4) \times 10^{-16}$	$(1.1 \pm 0.1) \times 10^{-15}$	$(6.3 \pm 0.6) \times 10^{-16}$	$(1.5 \pm 0.2) \times 10^{-15}$	$(2.4 \pm 0.2) \times 10^{-15}$
GJ 667C	6.8	0.46 <sup>w</sup>	$(5.2^{+0.9}_{-0.9}) \times 10^{-13}$	$5.3 \times 10^{-13}$	$(5.4 \pm 0.1) \times 10^{-14}$	$(5.1 \pm 0.4) \times 10^{-16}$	$(6.5 \pm 0.5) \times 10^{-16}$	$(7.6 \pm 0.6) \times 10^{-16}$	$(1.9 \pm 0.2) \times 10^{-15}$	$(3.0 \pm 0.2) \times 10^{-15}$
GJ 1214 <sup>d</sup>	14.6	0.21 <sup>z</sup>	$(1.3^{+1.4}_{-0.5}) \times 10^{-14}$	$1.4 \times 10^{-14}$	$(2.2 \pm 0.4) \times 10^{-15}$	$(8.0 \pm 2.8) \times 10^{-17}$	$(9.8 \pm 3.0) \times 10^{-17}$	$(4.3 \pm 3.4) \times 10^{-17}$	$(7.4 \pm 14.4) \times 10^{-17}$	$(5.2 \pm 1.2) \times 10^{-16}$
$\log_{10} T^f$			4.5 (core)		4.5 (core)	4.7	4.5	4.9	4.9	5.0

**Notes.**<sup>a</sup> Reconstructed value.<sup>b</sup>  $F(\text{EUV}) \Delta\lambda = 100\text{--}1170 \text{ \AA}$ .<sup>c</sup> We applied a correction factor of 1.3 to  $F(\text{Mg II})$  to correct for ISM attenuation (see Section 5.2).<sup>d</sup>  $F(\text{C II})$  is the only metal flux presented here that is not doublet integrated ( $\lambda 1335 \text{ \AA}$  only).<sup>e</sup> GJ 1214's  $\text{Ly}\alpha$  flux represents the scaled value determined from the  $\text{Ly}\alpha\text{--Mg II}$  surface flux relation.<sup>f</sup> All formation temperatures are from the CHIANTI database (Dere et al. 1997; Landi et al. 2013).<sup>g</sup> (t) Edvardsson et al. (1993), (u) Eker et al. (2015), (v) Bonfanti et al. (2015), (w) Kraus et al. (2011), (x) von Braun et al. (2014), (y) Zhou et al. (2014), (z) Berta et al. (2011).

$1.5 \times 10^{-5}$ , we assume that the effect of a fixed D/H ratio on the reconstructed Ly $\alpha$  fluxes is negligible.

### 3.3.3. Multiple Absorption Components

Since some of our sightlines are known to cross multiple interstellar clouds with different line-of-sight velocities (Redfield & Linsky 2008), we explore the effect of assuming a single ISM absorption component for all targets by adding a second absorption component to our fits for GJ 832, GJ 176, and HD 85512 ( $d = 4.9$  pc, 9.4 pc, and 11.2 pc).

Redfield & Linsky's (2008) dynamical model of the local ISM<sup>15</sup> has only one known interstellar cloud along GJ 832's line of sight (the Local Interstellar Cloud (LIC) with predicted radial velocity  $-11.3$  km s<sup>-1</sup> and predicted  $\log_{10} N(\text{H I}) = 16.6$ ). Two clouds are predicted along GJ 176's (the LIC with radial velocity  $23.7$  km s<sup>-1</sup> and  $\log_{10} N(\text{H I}) = 18.2$ , and the Hyades cloud with radial velocity  $12.9$  km s<sup>-1</sup>) sightline. Two clouds are also predicted along HD 85512's sightline (the G cloud with predicted radial velocity  $-0.1$  km s<sup>-1</sup>, and the Cet cloud with predicted radial velocity  $14.3$  km s<sup>-1</sup>). Predicted column densities for the LIC come from Redfield & Linsky's (2000) LIC model.<sup>16</sup> There are no H I column density predictions for any other local cloud.

For all three targets, we find the second ISM component to have  $\log_{10} N(\text{H I}) < 16$ , insignificant compared to the column densities of the first components (17.5–18.5). For GJ 832, we find the two ISM components have  $\log_{10} N(\text{H I}) = 18.2$  and  $14.4$  and  $V_{\text{H I}} = -17.1$  km s<sup>-1</sup> and  $-18.6$  km s<sup>-1</sup>. The reconstructed Ly $\alpha$  flux changed by  $< 2\%$ . For GJ 176, the two ISM components have  $\log_{10} N(\text{H I}) = 17.5$  and  $14.6$  and  $V_{\text{H I}} = 29.2$  km s<sup>-1</sup> and  $4.4$  km s<sup>-1</sup>. The reconstructed Ly $\alpha$  flux changed by  $< 2\%$ . For HD 85512, the two components have  $\log_{10} N(\text{H I}) = 18.3$  and  $15.7$  and  $V_{\text{H I}} = -2.8$  km s<sup>-1</sup> and  $-5.4$  km s<sup>-1</sup>. Although the quality of the fit did not change, the reconstructed Ly $\alpha$  flux decreased by 40%. The single-component and two-component fits are consistent within  $1\sigma$  error bars for all three stars.

We also simulated and fitted G140M and E140M observations of sightlines with 1–3 ISM absorption components and found that for the G140M observations, the MCMC-derived error bars are approximately  $1\sigma$  regardless of the number of ISM components along the sightline. For the E140M simulated observations with 2 and 3 ISM components, the MCMC-derived error bars may underestimate the true uncertainty. Depending on the ISM parameters and their velocity centroids compared to the intrinsic emission velocity centroids, we find over- and underestimates of 10%–30% of the true intrinsic flux.

We conclude that our low-resolution G140M and low signal-to-noise ratio (S/N) E140M Ly $\alpha$  profiles do not provide enough constraints on a second absorption component to justify its inclusion. We have shown in this section that for the G140M observations, there appears to be no systematic effect from excluding additional ISM components in the fits. For the E140M observations, the uncertainty introduced by ignoring a second or third ISM component could be as large as 30%. The two E140M-observed stars in the MUSCLES sample with multiple ISM components along the line of sight are HD 85512 and HD 97658.

### 3.3.4. Resolution Effects

We also explore the effects of resolution on our fits, because additional absorption features (multiple ISM components and astrospheric/heliospheric absorption) that are resolved in the E140M observations ( $R \sim 45,800$ ) will not be resolved in the G140M observations ( $R \sim 11,400$ ). We expect that the fitted parameters' uncertainties for high S/N G140M spectra should be larger than for high S/N E140M spectra, because resolving the D I absorption provides a tight constraint on the H I column density and velocity centroid, thereby reducing degeneracies in the model.

We choose three stars ( $\epsilon$  Eri, AU Mic, and 70 Oph A) with high S/N E140M observations from the *HST* archive and smooth the data to simulate a G140M observation. We smooth the Ly $\alpha$  profile to the approximate resolution of the G140M grating ( $R \sim 11,400$ ), and increase the error bar values to approximate the typical S/N per pixel of a G140M observation ( $\sim 20$ ). To quantify the absolute change in error bar size and test the stability of the best-fit parameters over many low-resolution “observations” with random noise, we add random noise (drawn from Gaussian distributions with varying standard deviation proportional to the error bars) and use our MCMC technique to fit the simulated observations 100 times.

For all three stars, we retrieve reconstructed Ly $\alpha$  fluxes  $\sim 18\%$  smaller than what we find from the original E140M spectra. The simulated G140M spectra have reconstructed profiles with larger FWHMs and smaller amplitudes for both the broad and narrow Gaussian components.  $A_{\text{narrow}}$  and  $A_{\text{broad}}$  on average decreased by 22% and 37%, respectively, and  $\text{FWHM}_{\text{narrow}}$  and  $\text{FWHM}_{\text{broad}}$  increased by 13% and 20%, respectively. Because the total intrinsic flux is proportional to  $A_{\text{narrow}} \times \text{FWHM}_{\text{narrow}} + A_{\text{broad}} \times \text{FWHM}_{\text{broad}}$ , these changes propagate to a 18% decrease in flux. H I column densities decreased by 7% on average, while Doppler  $b$  values increased by only a few percent. The marginalized error bars for the individual parameters and the reconstructed Ly $\alpha$  fluxes increase to levels consistent with the G140M observations ( $\sim 20\%$ ), although the stability of the reconstructed flux solution over the 100 trials is better ( $\sigma \sim 10\%$ ).

We list the specific findings for the three stars in Table 1. For reference,  $\epsilon$  Eri is a K dwarf with a single ISM velocity component and a detected astrosphere (the equivalent of the heliosphere around another star; Wood et al. 2005). AU Mic is an M dwarf with a single ISM velocity component and no detected astrospheric or heliospheric absorption, and 70 Oph A is a K dwarf with three known ISM velocity components and detected astrospheric and heliospheric absorption (Wood et al. 2005). For 70 Oph A, we use only one velocity component for the ISM absorption and make no attempt to treat the astrospheric and heliospheric absorption. At lower resolution, astrospheric absorption blends into the blueward wall of the H I absorption, and thus the fitted H I velocity centroid parameter shifts a few km s<sup>-1</sup> blueward and the Doppler  $b$  increases slightly.

We conclude that there likely is a 10%–20% systematic underestimate of the reconstructed fluxes for the G140M observations compared to the E140M observations, but note from our simulated observations that this underestimate is smaller than the  $1\sigma$  error bars about 68% of the time. We therefore do not correct our G140M results for this systematic error.

<sup>15</sup> <http://lism.wesleyan.edu/LISMdynamics.html>

<sup>16</sup> <http://lism.wesleyan.edu/ColoradoLIC.html>



We also call attention to the result that the derived  $N(\text{H I})$  and  $b$  values do not change appreciably when the spectral resolution is degraded to that corresponding to the G140M observations. This indicates that the  $N(\text{H I})$  and  $b$  values derived from analysis of the G140M spectra are accurate. We also note that from our simulated observations with multiple ISM components, we find that the fitted  $N(\text{H I})$  parameter approximately equals the sum of the true ISM components' column densities.

### 3.4. Comparison of Fitting Techniques

To verify the robustness of our MCMC optimizations, we also fit our model using MPFIT (Markwardt 2009) via the Python package `pyspeckit` (Ginsburg & Mirocha 2011). All reconstructed intrinsic Ly $\alpha$  fluxes agree within the MCMC results' error bars, although the individual parameters can vary more widely.

We also fit the M0 V star AU Mic with our MCMC technique to ascertain the agreement of our fitting technique and model with that of Wood et al. (2005) for a star not expected to show self-reversal of the Ly $\alpha$  line core. Our result shows excellent agreement with Wood et al.'s (2005) reconstructed Ly $\alpha$  flux and the ISM absorption parameters within a few percent. We find a total flux  $F(\text{Ly}\alpha) = (1.07 \pm 0.04) \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$ ,  $\log_{10} N(\text{H I}) = 18.35 \pm 0.01$ ,  $b = 14.0 \pm 0.2 \text{ km s}^{-1}$ , and  $V_{\text{H I}} = -24.0_{-0.3}^{+0.2} \text{ km s}^{-1}$ . Wood et al. (2005) found  $F(\text{Ly}\alpha) = 1.03 \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$ ,  $\log_{10} N(\text{H I}) = 18.356 \pm 0.002$ ,  $b = 13.74 \pm 0.08 \text{ km s}^{-1}$ , and  $V_{\text{H I}} = -20.83 \pm 0.09 \text{ km s}^{-1}$ .

### 3.5. Time Variability

The observed Ly $\alpha$  flux of our targets changes over time. We do not see any evidence for Ly $\alpha$  flaring in our STIS time-tag light curves, but multiple epochs of observations show some large changes. Figure 4 shows a comparison of the observed Ly $\alpha$  profiles for the seven targets with multiple observations and quantifies the changes in observed flux over time. GJ 832, GJ 667C, and  $\epsilon$  Eri exhibit significant changes in observed flux, and the other four stars are consistent with no change.

Between 2011 and 2012, the observed (not reconstructed) Ly $\alpha$  flux of GJ 832 increased by  $\sim 25\%$ , and by 2014, the observed flux decreased to  $\sim 15\%$  below the 2011 level. GJ 667C's observed profile appears broader in 2015 than it did in 2011, even after convolving the 2011 observed profile to the approximate resolution of the 2015 observation. The observed flux increased 25%, but we note that the 2011 observation has low S/N. For  $\epsilon$  Eri, the Ly $\alpha$  profiles (including the astrospheric absorption) are similar in the 1995 and 2015 observations, but they show a decrease of 14% in 2000.

GJ 876, GJ 436, and GJ 581 (assuming the apparent blueshift from 2012 to 2015 is an error in the wavelength calibration) exhibit no significant changes in their Ly $\alpha$  profile shapes and total fluxes. For GJ 1214, Ly $\alpha$  was not detected in 2011, and the observed flux ( $2.4 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$ ) is a  $1\sigma$  upper limit determined by integrating the spectrum over the wavelength range 1214.4–1216.9 Å. The 2015 observation (2.3 hr on-source with the  $52'' \times 0.''1$  slit) is more sensitive than the 2011 observation (2 hr on-source with the  $52'' \times 0.''05$  slit). The larger slit size and longer integration resulted in a  $5.5\text{-}\sigma$  detection of  $2.1 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$  for GJ 1214's observed flux, consistent with the 2011 upper limit.

The observed flux variations can be explained by activity cycles (year-to-decade timescales) and/or stellar rotation (day-to-week timescales). However, the time variations may be at least partially due to systematic effects. In comparing COS and STIS observations while creating the MUSCLES panchromatic spectra, Loyd et al. (2016) find systematic offsets between the two instruments in the tens of percent. The offsets can be explained by slit losses due to non-optimal centering of the target in the narrow STIS slit ( $52'' \times 0.''1$  for G140M and  $0.''2 \times 0.''06$  for E140M) in contrast to the minimal losses introduced by the larger COS aperture ( $2.''5$ ). Different observations may show different levels of accuracy in centering the target in the slit, despite acquisition peakup exposures, thereby producing artificial changes in the overall flux levels observed in the narrow-slit STIS G140M and E140M observations.

### 3.6. Extreme-UV Spectrum

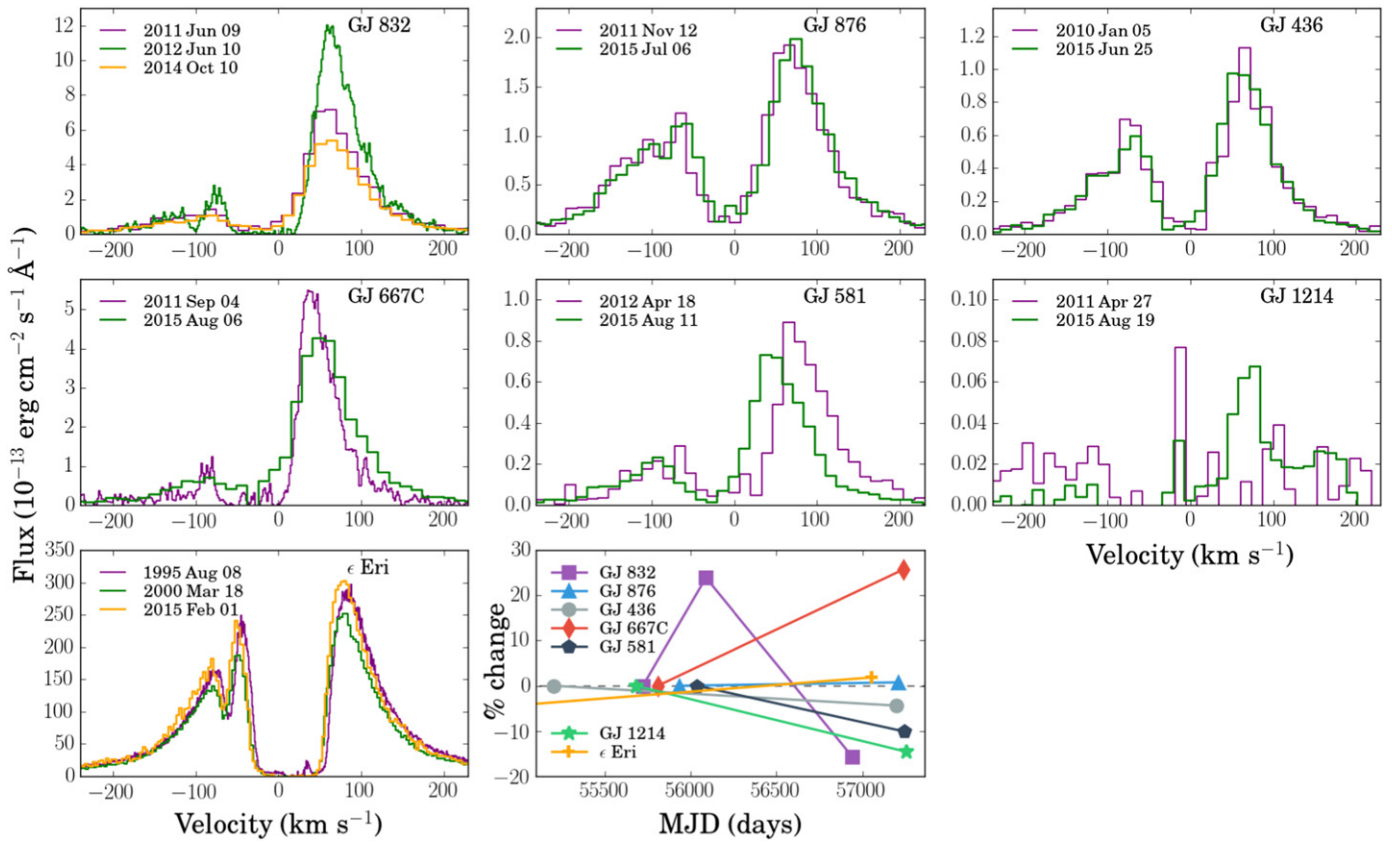
We compute the extreme-UV spectrum for each target in nine ( $\sim 100$  Å) bandpasses in the interval 100–1170 Å using empirically derived relations between total Ly $\alpha$  flux and extreme-UV flux presented by Linsky et al. (2014). Linsky et al. (2014) used solar models ( $\lambda < 2000$  Å) from Fontenla et al. (2013) to show that the ratio of extreme-UV flux to Ly $\alpha$  flux varies slowly with the Ly $\alpha$  flux. Using observations of F5 V–M5 V stars from *Far Ultraviolet Spectroscopic Explorer* (912–1170 Å) and *Extreme Ultraviolet Explorer* (100–400 Å), and intrinsic (reconstructed) Ly $\alpha$  fluxes from Wood et al. (2005) and Linsky et al. (2013), they established  $F(\Delta\lambda)/F(\text{Ly}\alpha)$  ratios, where  $\Delta\lambda$  represents one of the nine bandpasses in the interval 100–1170 Å. The estimated accuracy of these relations (the rms dispersions about the fit lines) is  $\sim 20\%$ . The extreme-UV fluxes ( $\Delta\lambda = 100\text{--}1170$  Å) are presented in Table 2, and the spectra are included in the High Level Science Products available on MAST. Example extreme-UV spectra for GJ 436, GJ 876, and HD 97658 are shown in Figure 5. In Section 4.3, we compare our extreme-UV flux estimate for GJ 436 to the estimate from Ehrenreich et al. (2015) and find agreement within 30%.

## 4. RESULTS

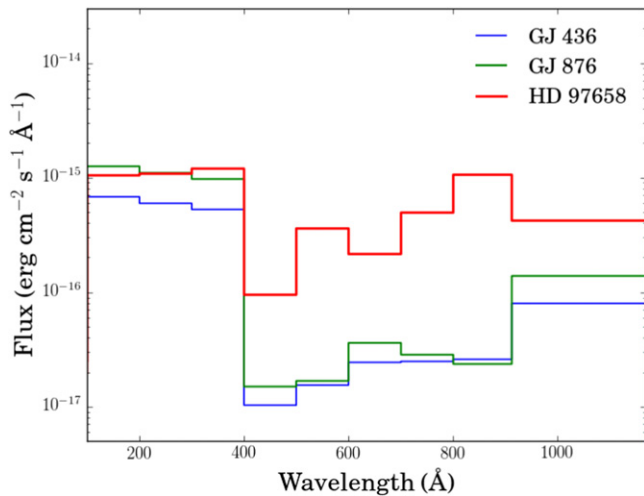
For each of the 11 target stars, the observed spectrum and its best-fit attenuated Ly $\alpha$  model are shown in Figure 1. See Tables 2 and 3 for the reconstructed Ly $\alpha$  fluxes and best-fit parameters. In the following subsections, we discuss the fitting and results in greater detail. We use the local ISM models of Redfield & Linsky (2000, 2008) to evaluate the plausibility of our ISM absorption best-fit parameters. For all 11 stars, we find H I velocity centroids that agree with the predictions of Redfield & Linsky (2008) within the MCMC uncertainties and STIS absolute velocity accuracies. The Redfield & Linsky (2000) model provides H I column density predictions only for the LIC, and our measured H I column densities do not always agree with the model predictions.

### 4.1. GJ 832

GJ 832 is a M1.5 V star at 5.0 pc observed with the low-resolution G140M grating. GJ 832's sightline is predicted to have only one ISM component, the LIC with radial velocity  $-11.3 \pm 1.3 \text{ km s}^{-1}$  (Redfield & Linsky 2008) and predicted  $\log_{10}$  H I column density 16.6 (Redfield & Linsky 2000).



**Figure 4.** First seven panels show multi-epoch Ly $\alpha$  profiles observed with *HST* STIS (the 1995 observation of  $\epsilon$  Eri was with GHRS). The last panel shows the percent changes in the observed Ly $\alpha$  fluxes. GJ 1214’s 0% value is based on an upper limit (Section 4.7). For  $\epsilon$  Eri, the year 2000 observed flux decreased by 14% relative to its 1995 value. The velocity shift between GJ 581’s two epochs is likely due to an error in the wavelength calibration.



**Figure 5.** Extreme-UV spectra of GJ 436, GJ 876, and HD 97658 from 100–1170 Å. These broadband spectra were constructed from the reconstructed stellar Ly $\alpha$  fluxes using the empirical relations presented in Linsky et al. (2014).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F$  (Ly $\alpha$ ) =  $(9.5 \pm 0.6) \times 10^{-13}$  erg cm $^{-2}$  s $^{-1}$ , and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.20 \pm 0.03$  and velocity centroid  $V_{\text{H I}} = -17.1^{+0.8}_{-0.7}$  km s $^{-1}$  (Tables 2 and 3). The column density is much higher than the predicted value for the LIC. The LIC may be clumpier than Redfield & Linsky’s (2000) model could account for, but it is possible there is an unknown

cloud component along this sightline. The best-fit velocity centroid is consistent with the prediction from Redfield & Linsky (2008) within G140M’s absolute wavelength accuracy (2.5–6.2 km s $^{-1}$ ).

GJ 832’s joint and marginalized probability distributions are shown in Figure 3 as the archetype for the high S/N (G140M or E140M) observations. The marginalized distributions are all singly peaked, and the Doppler  $b$  parameter’s marginalized distribution is the only one exhibiting significant asymmetry. Many joint distributions display degeneracies between parameters, and these are seen for all targets. See Section 3.2 for a description of the degeneracies.

#### 4.2. GJ 876

GJ 876 is a M5 V star at 4.7 pc observed with the low-resolution G140M grating. GJ 876’s sightline is predicted to have only one ISM component, the LIC with radial velocity  $-3.6 \pm 1.4$  km s $^{-1}$  (Redfield & Linsky 2008) and predicted  $\log_{10}$  H I column density 17.7 (Redfield & Linsky 2000).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F$  (Ly $\alpha$ ) =  $(3.9^{+0.5}_{-0.4}) \times 10^{-13}$  erg cm $^{-2}$  s $^{-1}$ , and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.03 \pm 0.04$  and velocity centroid  $V_{\text{H I}} = 2.8 \pm 0.8$  km s $^{-1}$  (Tables 2 and 3). The column density is twice as large as the predicted value for the LIC, and the velocity centroid is consistent within G140M’s absolute wavelength accuracy. As discussed in Section 4.1, this large best-fit column density could indicate substructure

**Table 3**  
Parameters from Ly $\alpha$  Reconstructions of MUSCLES Targets<sup>a</sup>

Target	$V_{\text{narrow}}$ (km s <sup>-1</sup> )	$A_{\text{narrow}}^b$ ( $\times 10^{-13}$ )	$\text{FWHM}_{\text{narrow}}$ (km s <sup>-1</sup> )	$V_{\text{broad}}$ (km s <sup>-1</sup> )	$A_{\text{broad}}^b$ ( $\times 10^{-13}$ )	$\text{FWHM}_{\text{broad}}$ (km s <sup>-1</sup> )	$\log_{10} N(\text{H I})$ (cm <sup>-2</sup> )	$b_{\text{H I}}$ (km s <sup>-1</sup> )	$V_{\text{H I}}$ (km s <sup>-1</sup> )	$\chi^2_\nu$
$\epsilon$ Eri	13.6 <sup>+0.4</sup> <sub>-0.4</sub>	937.4 <sup>+44.4</sup> <sub>-41.8</sub>	129.3 <sup>+1.6</sup> <sub>-1.6</sub>	11.7 <sup>+1.8</sup> <sub>-1.8</sub>	51.5 <sup>+2.8</sup> <sub>-2.6</sub>	400.1 <sup>+9.0</sup> <sub>-8.9</sub>	17.93 <sup>+0.02</sup> <sub>-0.02</sub>	12.1 <sup>+0.2</sup> <sub>-0.2</sub>	16.1 <sup>+0.4</sup> <sub>-0.4</sub>	1.9
HD 85512	-9.2 <sup>+3.3</sup> <sub>-4.0</sub>	22.7 <sup>+37.8</sup> <sub>-9.5</sub>	112.7 <sup>+16.8</sup> <sub>-23.1</sub>	-29.2 <sup>+11.3</sup> <sub>-31.5</sub>	1.3 <sup>+3.6</sup> <sub>-1.1</sub>	197.6 <sup>+81.1</sup> <sub>-32.3</sub>	18.43 <sup>+0.11</sup> <sub>-0.14</sub>	6.3 <sup>+6.6</sup> <sub>-1.4</sub>	-3.3 <sup>+0.7</sup> <sub>-0.8</sub>	0.8
HD 40307	17.3 <sup>+5.1</sup> <sub>-4.3</sub>	42.9 <sup>+50.5</sup> <sub>-21.5</sub>	102.2 <sup>+13.1</sup> <sub>-11.4</sub>	41.3 <sup>+29.8</sup> <sub>-12.3</sub>	2.1 <sup>+3.0</sup> <sub>-1.4</sub>	167.0 <sup>+32.4</sup> <sub>-21.4</sub>	18.60 <sup>+0.08</sup> <sub>-0.09</sub>	12.7 <sup>+1.5</sup> <sub>-2.6</sub>	9.4 <sup>+1.0</sup> <sub>-1.0</sub>	0.7
HD 97658	-1.1 <sup>+2.5</sup> <sub>-3.1</sub>	20.4 <sup>+23.1</sup> <sub>-10.3</sub>	99.0 <sup>+12.8</sup> <sub>-10.3</sub>	-37.6 <sup>+18.7</sup> <sub>-50.8</sub>	0.3 <sup>+0.3</sup> <sub>-0.2</sub>	252.6 <sup>+54.3</sup> <sub>-42.3</sub>	18.45 <sup>+0.09</sup> <sub>-0.14</sub>	6.0 <sup>+5.6</sup> <sub>-1.3</sub>	7.1 <sup>+0.9</sup> <sub>-1.0</sub>	0.8
GJ 832	11.4 <sup>+1.4</sup> <sub>-1.3</sub>	15.1 <sup>+1.1</sup> <sub>-1.0</sub>	133.4 <sup>+2.4</sup> <sub>-2.5</sub>	26.2 <sup>+4.2</sup> <sub>-4.0</sub>	0.49 <sup>+0.07</sup> <sub>-0.06</sub>	385.9 <sup>+17.5</sup> <sub>-16.5</sub>	18.20 <sup>+0.03</sup> <sub>-0.03</sub>	8.7 <sup>+1.0</sup> <sub>-1.8</sub>	-17.1 <sup>+0.8</sup> <sub>-0.7</sub>	1.3
GJ 876	7.8 <sup>+1.1</sup> <sub>-1.1</sub>	5.5 <sup>+0.7</sup> <sub>-0.6</sub>	146.8 <sup>+5.2</sup> <sub>-5.1</sub>	6.1 <sup>+5.9</sup> <sub>-5.7</sub>	0.23 <sup>+0.05</sup> <sub>-0.04</sub>	407.7 <sup>+31.6</sup> <sub>-26.8</sub>	18.03 <sup>+0.04</sup> <sub>-0.04</sub>	9.0 <sup>+1.3</sup> <sub>-2.7</sub>	2.8 <sup>+0.8</sup> <sub>-0.8</sub>	1.4
GJ 581	-2.6 <sup>+3.6</sup> <sub>-2.8</sub>	1.9 <sup>+0.5</sup> <sub>-0.4</sub>	133.0 <sup>+5.0</sup> <sub>-5.8</sub>	-38.8 <sup>+31.6</sup> <sub>-53.5</sub>	0.03 <sup>+0.02</sup> <sub>-0.01</sub>	509.4 <sup>+208.1</sup> <sub>-124.6</sub>	18.01 <sup>+0.12</sup> <sub>-0.17</sub>	12.0 <sup>+1.5</sup> <sub>-5.5</sub>	-24.1 <sup>+2.8</sup> <sub>-1.7</sub>	0.7
GJ 176	29.8 <sup>+0.6</sup> <sub>-0.6</sub>	4.6 <sup>+0.3</sup> <sub>-0.3</sub>	169.1 <sup>+3.2</sup> <sub>-3.2</sub>	31.4 <sup>+4.2</sup> <sub>-4.1</sub>	0.26 <sup>+0.03</sup> <sub>-0.03</sub>	498.2 <sup>+24.1</sup> <sub>-22.2</sub>	17.46 <sup>+0.08</sup> <sub>-0.10</sub>	15.9 <sup>+0.3</sup> <sub>-0.3</sub>	29.0 <sup>+0.4</sup> <sub>-0.4</sub>	1.2
GJ 436	2.6 <sup>+1.5</sup> <sub>-1.4</sub>	3.5 <sup>+0.6</sup> <sub>-0.5</sub>	126.7 <sup>+5.2</sup> <sub>-5.3</sub>	0.7 <sup>+8.2</sup> <sub>-7.6</sub>	0.12 <sup>+0.16</sup> <sub>-0.04</sub>	309.0 <sup>+46.1</sup> <sub>-36.6</sub>	18.04 <sup>+0.06</sup> <sub>-0.06</sub>	8.6 <sup>+1.4</sup> <sub>-2.4</sub>	-4.1 <sup>+1.3</sup> <sub>-1.2</sub>	0.9
GJ 667C	12.7 <sup>+4.5</sup> <sub>-3.1</sub>	8.7 <sup>+1.7</sup> <sub>-1.6</sub>	126.7 <sup>+3.7</sup> <sub>-3.9</sub>	22.1 <sup>+7.0</sup> <sub>-6.9</sub>	0.28 <sup>+0.06</sup> <sub>-0.05</sub>	383.0 <sup>+30.5</sup> <sub>-26.9</sub>	17.98 <sup>+0.12</sup> <sub>-0.18</sub>	10.6 <sup>+2.0</sup> <sub>-5.0</sub>	-22.5 <sup>+2.3</sup> <sub>-1.6</sub>	1.1
GJ 1214 <sup>c</sup>	32.4 <sup>+31.7</sup> <sub>-21.2</sub>	0.18 <sup>+0.20</sup> <sub>-0.06</sub>	166.0 <sup>+27.1</sup> <sub>-27.3</sub>	...	...	...	18.06 <sup>+0.90</sup> <sub>-0.42</sub>	14.9 <sup>+3.8</sup> <sub>-6.2</sub>	-26.4 <sup>+2.2</sup> <sub>-82.0</sub>	...

**Notes.**

<sup>a</sup> D/H fixed at  $1.5 \times 10^{-5}$ .

<sup>b</sup> erg cm<sup>-2</sup> s<sup>-1</sup> Å<sup>-1</sup>.

<sup>c</sup> GJ 1214's amplitude represents the scaled value determined from the Ly $\alpha$ -Mg II relation.

within the LIC or an unknown cloud component along this sightline.

GJ 876's joint and marginalized distributions are singly peaked and similar in shape and symmetry to Figure 3, including the asymmetry in the marginalized distribution for the Doppler  $b$  value.

#### 4.3. GJ 436

GJ 436 is a M3.5 V star at 10.1 pc observed with the low-resolution G140M grating. According to the Redfield & Linsky (2008) model, GJ 436's sightline does not pass directly through any local cloud, but is close to or inside the edges of three clouds (the LIC, Leo, and NGP). Redfield & Linsky (2000) predicted a  $\log_{10} N(\text{H I})$  column density of 16.7 associated with the LIC.

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (2.1 \pm 0.3) \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup>, and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.04 \pm 0.06$  and velocity centroid  $V_{\text{H I}} = -4.1^{+1.3}_{-1.2}$  km s<sup>-1</sup> (Tables 2 and 3). The column density is consistent with other 10 pc distant stars reported in Wood et al. (2005), indicating the presence of significant interstellar material.

GJ 436's joint and marginalized distributions of the individual parameters are singly peaked and similar in shape and symmetry to Figure 3, including the asymmetry in the marginalized distribution for the Doppler  $b$  value.

Bourrier et al. (2015) observed GJ 436 over two *HST* visits and performed Ly $\alpha$  reconstructions assuming a single Voigt profile (with no self-reversal) as the intrinsic Ly $\alpha$  shape. Our  $F(\text{Ly}\alpha)$  and  $\log_{10} N(\text{H I})$  values are in good agreement with theirs:  $F(\text{Ly}\alpha) = 2.01 \times 10^{-13}$  and  $2.05 \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup> and  $\log_{10} N(\text{H I}) = 18.01 \pm 0.07$  and  $17.90^{+0.15}_{-0.19}$ .

We have also compared our extreme-UV flux computed using our reconstructed Ly $\alpha$  flux (Section 3.6) with the extreme-UV flux computed by Ehrenreich et al. (2015) using the X-ray flux from 5–100 Å, via relations provided by Sanz-Forcada et al. (2011). We find that our extreme-UV fluxes are in satisfactory agreement. Ehrenreich et al. (2015) found  $F(\text{EUV}) = 2.34 \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup> over

$\Delta\lambda = 124\text{--}912$  Å. Integrating our reconstructed extreme-UV over the same wavelength range, we find  $F(\text{EUV})^{17} = 1.7 \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup>, 27% less than what Ehrenreich et al. (2015) found.

#### 4.4. GJ 581

GJ 581 is a M5 V star at 6.2 pc observed with the low-resolution G140M grating. GJ 581's sightline is predicted to have two ISM components: the G cloud with a radial velocity of  $-27.6 \pm 1.2$  km s<sup>-1</sup> and the Gem cloud with radial velocity  $-23.8 \pm 1.0$  km s<sup>-1</sup> (Redfield & Linsky 2008).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (1.1^{+0.2}_{-0.3}) \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup>, and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.01^{+0.12}_{-0.17}$  and velocity centroid  $V_{\text{H I}} = -24.1^{+2.8}_{-1.7}$  km s<sup>-1</sup> (Tables 2 and 3). The velocity centroid is consistent with the predicted radial velocities within uncertainties for both clouds.

GJ 581's joint and marginalized distributions are singly peaked except for the Doppler  $b$  value and exhibit a quality of shape and symmetry halfway between Figure 3. Many of the parameters, while singly peaked, exhibit significant asymmetry in their marginalized distributions. This is reflected in the uncertainties listed in Tables 2 and 3.

#### 4.5. GJ 667C

GJ 667C is a M1.5 V star at 6.8 pc observed with the low-resolution G140M grating. GJ 667C's sightline is predicted to have one ISM component, the G cloud with a radial velocity of  $-27.3 \pm 1.2$  km s<sup>-1</sup> (Redfield & Linsky 2008).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (5.2 \pm 0.9) \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup>, and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 17.98^{+0.12}_{-0.18}$  and velocity centroid  $V_{\text{H I}} = -22.5^{+2.3}_{-1.6}$  km s<sup>-1</sup> (Tables 2 and 3). The velocity centroid is consistent within G140M's absolute wavelength accuracy of the predicted radial velocity.

<sup>17</sup> The  $F(\text{EUV}) = 2.1 \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup> value reported in Table 2 is computed over  $\Delta\lambda = 100\text{--}1170$  Å.

GJ 667C’s joint and marginalized distributions are singly peaked except for the Doppler  $b$  value and exhibit at least a slight asymmetry for most of the parameters and significant asymmetry for the Doppler  $b$  value. These asymmetries are reflected in the uncertainties listed in Tables 2 and 3.

#### 4.6. GJ 176

GJ 176 is a M2.5 V star at 9.3 pc observed with the low-resolution G140M grating. GJ 176’s sightline is predicted to have two ISM components: the LIC with radial velocity  $23.7 \pm 0.9 \text{ km s}^{-1}$  and the Hyades cloud with radial velocity  $12.9 \pm 1.1 \text{ km s}^{-1}$  (Redfield & Linsky 2008).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (3.9 \pm 0.2) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ , and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 17.46_{-0.10}^{+0.08}$  and velocity centroid  $V_{\text{H I}} = 29.0 \pm 0.4 \text{ km s}^{-1}$  (Tables 2 and 3). The velocity centroid is consistent with the predicted radial velocities of the LIC and Hyades cloud within G140M’s absolute wavelength accuracy. Redfield & Linsky (2000) predicted that the LIC along GJ 176’s sightline has a  $\log_{10}$  H I column density 18.2, a factor of four larger than our measurement. GJ 176’s H I column density is among the lowest column densities measured for any star and is extremely low for a star at 9.4 pc. However, the best fit has an anomalously large Doppler  $b$  value,  $15.9 \pm 0.3 \text{ km s}^{-1}$ . Forcing the  $b$  value to a more probable  $12 \text{ km s}^{-1}$  increases the H I column density by 30%, but changes the reconstructed Ly $\alpha$  flux by only a few percent.

While we are skeptical of the fitted ISM parameter values, we have reproduced this solution with MPFIT and a gridsearch. GJ 176 is also the highest S/N observation of the MUSCLES sample, and its singly peaked, symmetric joint and marginalized distributions indicate a good fit. If GJ 176 is among the sightlines with the lowest column densities, it would make an excellent target for future direct extreme-UV observations.

#### 4.7. GJ 1214

GJ 1214 is a M4.5 V star at 14.6 pc observed with the low-resolution G140M grating. GJ 1214’s sightline is predicted to have two ISM components: the LIC with radial velocity  $-22.3 \pm 1.0 \text{ km s}^{-1}$  and the Mic cloud with radial velocity  $-26.6 \pm 1.0 \text{ km s}^{-1}$  (Redfield & Linsky 2008).

We estimate a total flux  $F(\text{Ly}\alpha) = (1.3_{-0.5}^{+1.4}) \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$ , and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.06_{-0.42}^{+0.90}$  and velocity centroid  $V_{\text{H I}} = -26.4_{-82.0}^{+22.2} \text{ km s}^{-1}$  (Tables 2 and 3). None of the parameters is well constrained. The predicted  $\log_{10}$  H I column density associated with the LIC along this sightline is 16.3, so the Mic cloud probably provides the dominant absorption. The best-fit H I velocity centroid agrees within 1% of the predicted radial velocity of the Mic cloud, although its uncertainty is much larger.

In the MUSCLES pilot program (France et al. 2013), Ly $\alpha$  emission was not detected from GJ 1214 with a  $1\sigma$  upper limit  $2.4 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$ . We now detect Ly $\alpha$  at 5.5- $\sigma$  significance ( $2.1 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$ ), but the S/N is so low that we include only a single emission component in the model. The fit is not well constrained, with the marginalized distributions not peaking for the column density and Doppler  $b$  parameters over physically reasonable ranges ( $17.5 \leq \log_{10} N(\text{H I}) \leq 20$ ;  $1 \leq b \leq 20 \text{ km s}^{-1}$ ). The other marginalized

distributions, while singly peaked, exhibit large asymmetries as reflected in the reported uncertainties (Table 3).

Because the fit quality is poor, we attempt to further constrain GJ 1214’s intrinsic Ly $\alpha$  emission using an empirical Ly $\alpha$ –Mg II surface flux relation (see Section 5.2). Because these two transitions originate from gas with similar temperature and opacity, their fluxes are tightly correlated. Comparing the MCMC reconstructed flux to the flux expected from the empirical Ly $\alpha$ –Mg II relation, we are likely underestimating GJ 1214’s intrinsic Ly $\alpha$  flux by a factor of 2.4. The best-fit line in Figure 1 (pink line) does appear to be an underestimate, and the orange line shows the scaled profile (the amplitude parameter has been increased by a factor of 2.4). The scaled profile’s amplitude appears too large to match the observed profile, but the observed profile’s flux is likely underestimated due to a possible oversubtraction of airglow. Thus, we present the scaled flux and individual parameters (only the amplitude has been modified (2.4 $\times$ ) from the original MCMC fit) in Tables 2 and 3.

#### 4.8. HD 85512

HD 85512 is a K6 V star at 11.2 pc observed with the medium-resolution E140M grating. HD 85512’s sightline is predicted to have two ISM components: the G cloud with a radial velocity of  $-0.1 \pm 1.0 \text{ km s}^{-1}$  and the Cet cloud with radial velocity  $14.3 \pm 0.7 \text{ km s}^{-1}$  (Redfield & Linsky 2008).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (1.2_{-0.5}^{+1.8}) \times 10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$ , and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.43_{-0.14}^{+0.11}$  and velocity centroid  $V_{\text{H I}} = -3.3_{-0.8}^{+0.7} \text{ km s}^{-1}$  (Tables 2 and 3). The velocity centroid is consistent with the G cloud, but not the Cet cloud. The Cet cloud may not significantly contribute to the absorption along this sightline.

HD 85512’s joint and marginalized distributions are shown in Figure 4 as the archetype of the E140M low-S/N observations, although it is actually the worst case. The distributions are doubly peaked for half the parameters, and all but the H I velocity centroid exhibit significant asymmetry. The marginalized distribution of the broad component’s velocity centroid is flat for values less than  $40 \text{ km s}^{-1}$ , so not to skew the best-fit value significantly, we limit the short wavelength end of the parameter range  $50 \text{ km s}^{-1}$  further than for the narrow component’s velocity centroid. We also limit the broad component’s amplitude to two orders of magnitude fainter than the narrow component’s range, because the MCMC cannot rule out a broad component consistent with zero flux. These requirements we had to place on the broad emission component’s parameters indicate that the broad component is only marginally detected.

#### 4.9. HD 97658

HD 97658 is a K1 V star at 21.1 pc observed with the medium-resolution E140M grating. HD 97658’s sightline is predicted to have three ISM components: the LIC with predicted radial velocity of  $2.7 \pm 1.4 \text{ km s}^{-1}$ , the Leo cloud with predicted radial velocity  $4.7 \pm 0.8 \text{ km s}^{-1}$ , and the NGP with predicted radial velocity  $9.5 \pm 0.8 \text{ km s}^{-1}$  (Redfield & Linsky 2008). Redfield & Linsky (2000) predicted a  $\log_{10}$  H I column density of 16.8 from the LIC along this sightline.

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (9.1_{-4.5}^{+9.9}) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ , and the best-fit H I

absorption has  $\log_{10} N(\text{H I}) = 18.45^{+0.09}_{-0.14}$  and velocity centroid  $V_{\text{H I}} = 7.1^{+0.9}_{-1.0}$  km s<sup>-1</sup> (Tables 2 and 3). The velocity centroid is consistent within E140M’s absolute wavelength accuracy (0.7–1.6 km s<sup>-1</sup>) for the Leo and NGP clouds. The LIC is probably not the dominant source of absorption along this line of sight.

HD 97658’s joint and marginalized distributions are singly peaked, except for the Doppler  $b$  value, and exhibit at least slight asymmetry for most of the parameters. These asymmetries are reflected in the uncertainties listed in Tables 2 and 3. The low S/N of this observation results in a large degeneracy between a low  $b$  value (6 km s<sup>-1</sup>), high column density (18.45) solution and a high  $b$  value (12 km s<sup>-1</sup>), low column density (18.2) solution. The low  $b$  value, high column density fits the observed profile best, but the 12 km s<sup>-1</sup> is more realistic. Restricting the Doppler  $b$  value to be  $>9$  km s<sup>-1</sup> results in  $F(\text{Ly}\alpha) = (4.8^{+3.7}_{-1.7}) \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup>, which is at the lower end of the original  $1\sigma$  error bars.

#### 4.10. HD 40307

HD 40307 is a K2.5 V star at 13.0 pc observed with the medium-resolution E140M grating. HD 40307’s sightline is predicted to have only one ISM component, the Blue cloud with a radial velocity of  $7.6 \pm 1.1$  km s<sup>-1</sup> (Redfield & Linsky 2008).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (2.0^{+2.2}_{-0.9}) \times 10^{-12}$  erg cm<sup>-2</sup> s<sup>-1</sup>, and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 18.60^{+0.08}_{-0.09}$  and velocity centroid  $V_{\text{H I}} = 9.4 \pm 1.0$  km s<sup>-1</sup> (Tables 2 and 3). The velocity centroid is consistent within uncertainties with the predicted radial velocity.

HD 40307’s joint and marginalized distributions are all singly peaked, although the Doppler  $b$  value features a broad hump on its lower tail and the broad flux component’s velocity centroid distribution becomes flat and non-zero at large velocities. This flat part of the distribution could significantly skew the median value away from the peak of the distribution, so we restricted the broad velocity centroid’s parameter range extend 70 km s<sup>-1</sup> farther than the narrow component’s range, because it seemed unlikely that the two flux components could be offset by such a large velocity. The necessity of this restriction indicates that the broad emission component is only marginally detected.

#### 4.11. $\epsilon$ Eri

$\epsilon$  Eri is a K2 V star at 3.2 pc observed with the medium-resolution E140M grating and is our highest S/N observation with that grating.  $\epsilon$  Eri’s sightline is measured to have only one ISM component, the LIC with radial velocity  $20.9 \pm 1.8$  km s<sup>-1</sup> (Redfield & Linsky 2008) and column density  $17.88 \pm 0.07$  (Dring et al. 1997).

The best-fit reconstructed Ly $\alpha$  profile has a total flux  $F(\text{Ly}\alpha) = (6.1^{+0.3}_{-0.2}) \times 10^{-11}$  erg cm<sup>-2</sup> s<sup>-1</sup>, and the best-fit H I absorption has  $\log_{10} N(\text{H I}) = 17.93 \pm 0.02$  and velocity centroid  $V_{\text{H I}} = 16.1 \pm 0.4$  km s<sup>-1</sup> (Tables 2 and 3). The velocity centroid is not quite consistent with the measured radial velocity, but the column density agrees with the measurement from Dring et al. (1997).  $\epsilon$  Eri’s joint and marginalized distributions are all singly peaked and none display asymmetry.

**Table 4**  
Ly $\alpha$  Broad Component Strengths Compared to Total Intrinsic Ly $\alpha$  Flux

Target	$F_{\text{broad}}/F_{\text{total}}$ <sup>a</sup> (%)	$F_{\text{broad wings}}/F_{\text{total}}$ <sup>b</sup> (%)
$\epsilon$ Eri	15	... <sup>c</sup>
HD 85512	9	0.9
HD 40307	7	0.7
HD 97658	4	1.1
GJ 832	9	1.3
GJ 876	10	2.2
GJ 581	7	0.9
GJ 176	14	2.7
GJ 436	8	1.0
GJ 667C	9	0.6
GJ 1214	0 <sup>d</sup>	0

#### Notes.

<sup>a</sup> Percentage of the total reconstructed Ly $\alpha$  flux ( $F_{\text{total}}$ ) contained in the broad Gaussian component ( $F_{\text{broad}}$ ).

<sup>b</sup> Percentage of the total reconstructed Ly $\alpha$  flux ( $F_{\text{total}}$ ) contained in the broad wings ( $F_{\text{broad wings}}$  is determined over  $\Delta\lambda = \pm(300\text{--}1200$  km s<sup>-1</sup>) from the COS G130M data).

<sup>c</sup>  $\epsilon$  Eri does not have COS G130M data.

<sup>d</sup> GJ 1214 was fit with only 1 Gaussian component.

$\epsilon$  Eri is the only star in our sample to show the presence of an astrosphere, as observed by Wood et al. (2005). To fit the attenuated profile, we mask the blue side of the H I absorption where the astrospheric absorption is present (Figure 1) because the extra absorption skews the H I velocity centroid blueward and increases the Doppler  $b$  value. The detected D I absorption eliminates the possibility of this blueward velocity centroid.

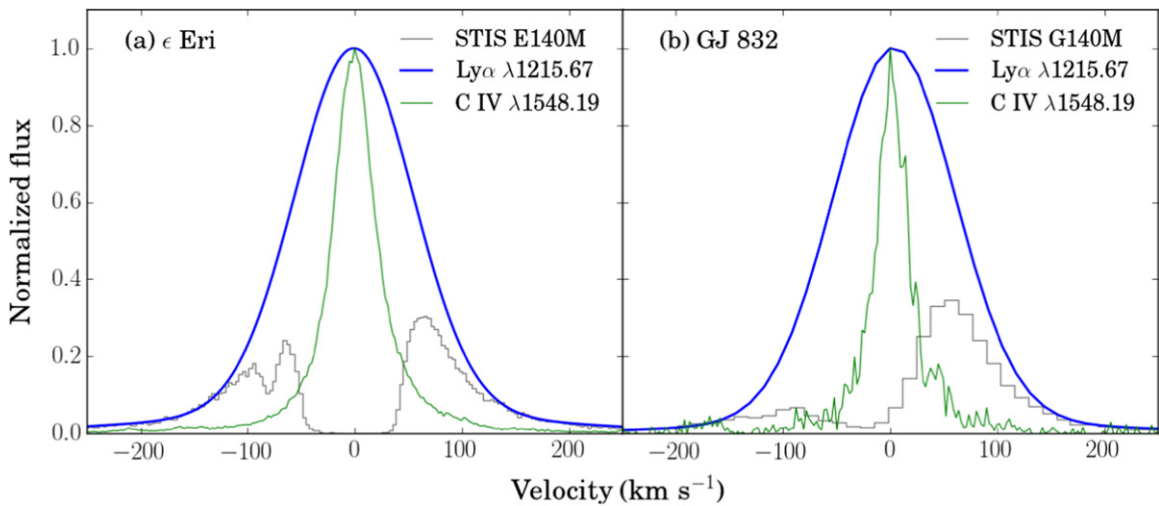
## 5. DISCUSSION

### 5.1. The Ly $\alpha$ Line Profile

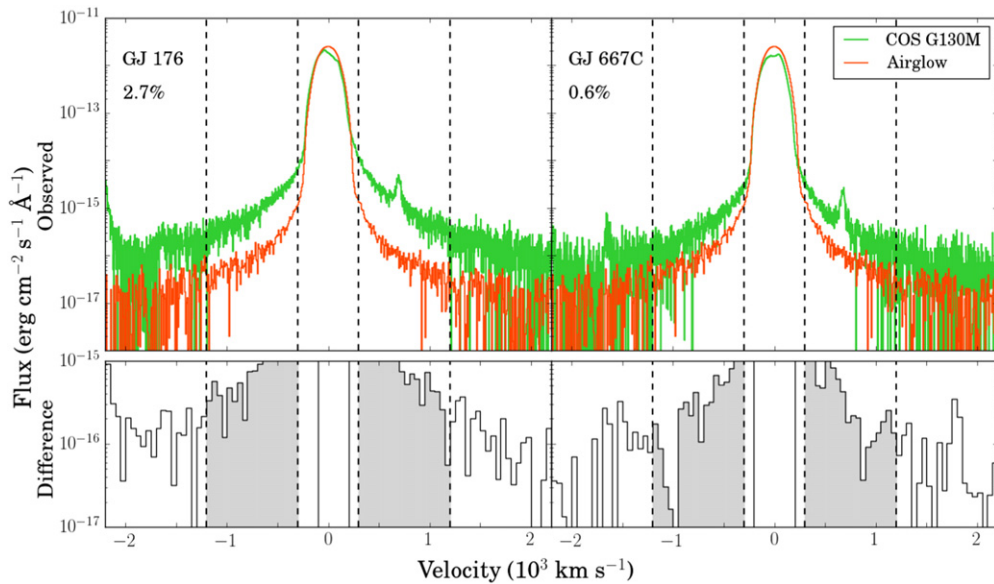
The Ly $\alpha$  emission lines observed by STIS have broad wings that are well fit by a narrow and a broad Gaussian. Excluding our targets with poor S/N where the broad component is only marginally detected (HD 85512, HD 97658, HD 40307) or not at all (GJ 1214), the average and standard deviation of the narrow component FWHMs is  $138 \pm 14$  km s<sup>-1</sup>, and  $413 \pm 64$  km s<sup>-1</sup> for the broad components. The broad flux component typically contributes only 10% to the total intrinsic flux (Table 4).

We compare the intrinsic Ly $\alpha$  profiles with the C IV profiles ( $\lambda 1550$ ) and find that Ly $\alpha$  is  $\sim 3$  times broader (Figure 6). C IV has a higher formation temperature ( $\log_{10} T = 5.0$ ) than the Ly $\alpha$  core ( $\log_{10} T = 4.5$ ) and forms only in the transition region while Ly $\alpha$  forms in both the chromosphere and transition region. The abundance of neutral hydrogen is also much higher than triply ionized carbon, and the measured ratio between the two members of the C IV doublet is  $\sim 2$ . This confirms that C IV is optically thin, while photons must scatter out of the optically thick Ly $\alpha$  line core to escape the stellar atmosphere.

The COS G130M Ly $\alpha$  spectra, while overwhelmed by airglow in the inner  $\pm 300$  km s<sup>-1</sup> of the line core, are more sensitive to the broad wings of the Ly $\alpha$  profile than the STIS spectra are. We detect broad Ly $\alpha$  wings out to  $\pm 1200$  km s<sup>-1</sup>. This flux probably originates from photons diffusing out of the Ly $\alpha$  line core. To estimate how much of the total flux is in the



**Figure 6.** Thin gray histogram shows the observed Ly $\alpha$  spectrum, the thick blue line shows the reconstructed Ly $\alpha$  profile, and the thin green line shows the observed C IV ( $\lambda 1548$ ) profile, plotted on a common velocity scale. We have defined 0 km s $^{-1}$  to be where each profile peaks. (a)  $\epsilon$  Eri. (b) GJ 832.

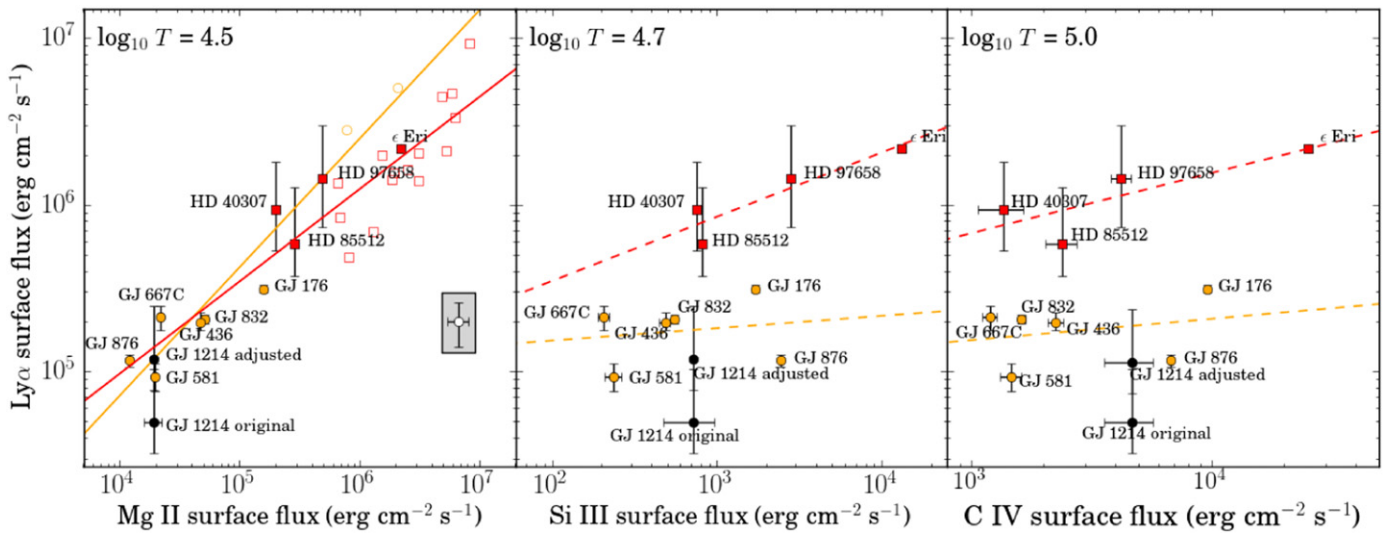


**Figure 7.** Broad Ly $\alpha$  wings of two of our targets, GJ 176 and GJ 667C. The “observed” flux panels show the Ly $\alpha$  profiles observed with COS G130M (green) and the template COS G130M airglow spectrum (orange-red). The “difference” flux panels show the COS G130M spectrum with the COS airglow template subtracted and the O V emission line at +650 km s $^{-1}$  removed. The dashed vertical lines and gray shading show the velocity ranges ( $-1200$  to  $-300$  km s $^{-1}$  and  $+300$  to  $+1200$  km s $^{-1}$ ) over which the “difference” flux was integrated. The percentage of the total Ly $\alpha$  flux contained in the gray-shaded regions is listed in each panel.

broad wings (beyond  $\pm 300$  km s $^{-1}$ ), we subtract a template COS G130M airglow spectrum from the observed G130M spectrum, interpolate over the O V emission line at +650 km s $^{-1}$ , and integrate the flux between  $\pm(300-1200)$  km s $^{-1}$  (Figure 7). We find percentages of the total Ly $\alpha$  flux in the wings as high as 2.7% (GJ 176) and as low as 0% (GJ 1214) with an average of 1.1% (Table 4). However, the percentage flux in the broad wings depends on the STIS observation’s S/N. As the S/N increases, the MCMC finds larger FWHM values for both the narrow and broad flux components, because the decrease in the noise floor allows for better sampling of the profile at large velocities. The S/N dependence indicates that the percentages listed in Figure 7 are lower limits, thus we estimate that  $>1\%$  of the total Ly $\alpha$  flux comes from the  $\pm(300-1200)$  km s $^{-1}$  wings of the MUSCLES targets.

## 5.2. Comparison with Other Chromospheric and Transition Region Lines

Correlations between reconstructed Ly $\alpha$  fluxes and other stellar emission lines will provide future studies with an alternative to estimate the intrinsic Ly $\alpha$  flux and can aid the development of low-mass stellar atmospheric models. We compare Ly $\alpha$  to metal ions of chromospheric or transition region origin whose observed fluxes and formation temperatures are listed in Table 2. Many of the observed lines are multiplets, and for all but C II (whose  $\lambda 1334$  line is attenuated by the ISM), we derive the total flux by integrating over all multiplet members. Both Mg II ( $\lambda\lambda 2796, 2803$ ) lines lose  $\sim 30\%$  of their flux to interstellar attenuation (France et al. 2013), assuming a typical  $\log_{10} N(\text{Mg II}) \sim 13$  for stars within 20 pc (Redfield & Linsky 2002), so we adjust the Mg II flux values by +30% (Table 2). The COS and STIS G230L



**Figure 8.** Left: Ly $\alpha$  surface flux vs. Mg II surface flux (corrected for ISM absorption). The M dwarfs are shown with orange circles and the K dwarfs in red squares (filled for the MUSCLES sample, open for the Wood et al. 2005 sample). GJ 1214’s original and adjusted fluxes are shown with black circles because GJ 1214 was excluded from the fits. The best-fit lines to the MUSCLES and Wood et al. (2005) data are shown in orange for the M dwarfs and red for the K dwarfs. For visual clarity, we show the size of the Wood et al. (2005) error bars in the gray-shaded box. Middle: Ly $\alpha$  surface flux vs. Si III surface flux. Right: Ly $\alpha$  surface flux vs. C IV surface flux. For the middle and right panels, the lines of best fit are shown as dashed lines because they are not statistically significant correlations (see Table 5). The formation temperatures of the metal ions are shown in the top left corner of each panel (Dere et al. 1997; Landi et al. 2013).

**Table 5**  
Parameters for Fits in Figures 8 and 9

Equation	Spectral Type	$\alpha^a$	$\beta^a$	$p$ -value <sup>b</sup>	$\sigma^c$
$\log_{10} F(\text{Ly}\alpha) = \alpha \log_{10} F(\text{Mg II}) + \beta$	M V	$0.77 \pm 0.10$	$1.77 \pm 0.52$	$1.0 \times 10^{-4}$	0.029
	K V	$0.55 \pm 0.11$	$2.77 \pm 0.68$	$4.3 \times 10^{-5}$	0.031
$\log_{10} F(\text{Ly}\alpha) = \alpha \log_{10} F(\text{Si III}) + \beta$	M V	$0.07 \pm 0.31$	$5.04 \pm 0.87$	0.75	0.033
	K V	0.39	4.77	0.086	0.015
$\log_{10} F(\text{Ly}\alpha) = \alpha \log_{10} F(\text{C IV}) + \beta$	M V	$0.13 \pm 0.35$	$4.80 \pm 1.21$	0.63	0.033
	K V	0.36	4.74	0.19	0.021
$\log_{10} F(\text{Ly}\alpha) = \alpha \log_{10} P + \beta$	M V	$-0.86 \pm 0.16$	$6.89 \pm 0.24$	$3.3 \times 10^{-4}$	0.037
	K V	$-0.54 \pm 0.05$	$6.79 \pm 0.05$	$3.2 \times 10^{-10}$	0.019

#### Notes.

<sup>a</sup> Uncertainties were determined by the fit’s covariance matrix. For the K-dwarf Si III and C IV relations, four data points are insufficient to calculate a covariance matrix.

<sup>b</sup> The  $p$ -value is the probability of finding the best-fit line assuming the null hypothesis (no correlation) is true.

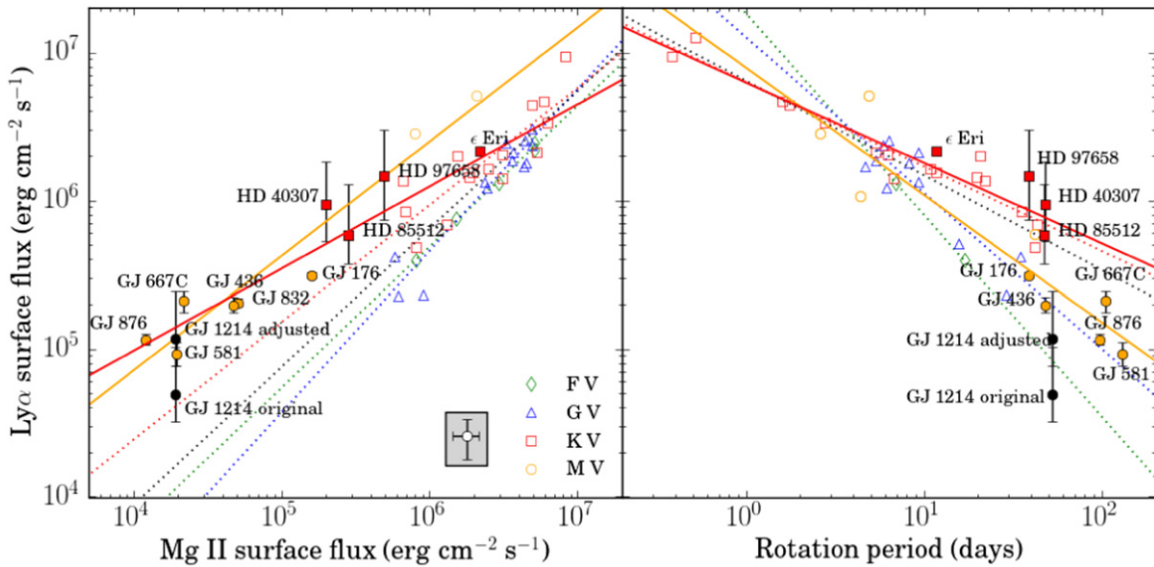
<sup>c</sup> The standard deviation of the data points about the best-fit line, normalized by the fit.

spectra of Mg II do not provide sufficient resolution for profile reconstruction. We do not adjust the  $F(\text{Mg II})$  correction factor based on the  $N(\text{H I})$  value measured for each sightline, because the  $\text{Mg}^+/\text{H}$  is not fixed for the local ISM. It is sensitive to abundance variations as well as local temperature and ionization conditions. The list of compared ions is: C II ( $\lambda 1335$ ), Mg II ( $\lambda \lambda 2796, 2803$ ), Si III ( $\lambda \lambda 1206, 1206$ ), Si IV ( $\lambda \lambda 1394, 1403$ ), He II ( $\lambda 1640$ ), and C IV ( $\lambda \lambda 1548, 1551$ ).

We expect that the surface fluxes of species with similar formation temperatures will correlate well (France et al. 2013; Linsky et al. 2013). Surface fluxes are calculated using the distances and stellar radii listed in Table 2. Figure 8 compares the surface fluxes of Ly $\alpha$  and three ions: Mg II (line center  $\log_{10} T = 4.5$ ), Si III ( $\log_{10} T = 4.7$ ), and C IV ( $\log_{10} T = 5.0$ ). In the Mg II panel, we include the M and K-dwarf samples from Wood et al. (2005) in the fit. We fit a line to the data for both spectral types (best-fit equations printed Table 5), but exclude GJ 1214 from each fit, because its reconstructed Ly $\alpha$  flux is highly uncertain. Because we are mixing data sets that have different error analyses, we use uniform weighting for all the

data points during the fit. As expected, we find positive correlations with lines whose formation temperature is closer to Ly $\alpha$ ’s line center  $\log_{10} T = 4.5$ . Some scatter is expected due to time variability (the ions’ spectra were taken close in time to Ly $\alpha$ ’s, but are not simultaneous) and slight differences or ranges in formation temperature.

Here we evaluate the statistical significance of trends revealed in Figure 8. First, the Ly $\alpha$ –Mg II correlations are the only ones with statistical significance ( $p$ -value  $< 0.05$ ; Table 5). The slopes of the K- and M-dwarf Ly $\alpha$ –Mg II best-fit lines are statistically consistent with each other, as are the K- and M-dwarf slopes for the Si III and C IV relations shown in Figure 8. Therefore, we cannot conclude that our sample shows a steeper Ly $\alpha$ –Mg II correlation for the M dwarfs than the K dwarfs. Figure 8 shows that the M-dwarf best-fit lines for Si III and C IV have shallower slopes compared to the M-dwarf Ly $\alpha$ –Mg II correlation, but these fits are not statistically significant, probably due to our small sample size (4 K dwarfs and 6 M dwarfs; GJ 1214 was excluded from fitting). We do not show the fitting results for C II, Si IV, or He II, but the fit qualities



**Figure 9.** Left:  $\text{Ly}\alpha$  surface flux vs.  $\text{Mg II}$  surface flux (corrected for ISM absorption). The M dwarfs are shown with orange circles and the K dwarfs in red squares (filled for the MUSCLES sample, open for the Wood et al. 2005 sample). GJ 1214’s original and adjusted fluxes are shown with black circles because GJ 1214 was excluded from the least-squares fits. The best-fit lines to the MUSCLES and Wood et al. (2005) data are shown in orange for the M dwarfs and red for the K dwarfs. The black dotted line is the least-squares fit to all of the Wood et al. (2005) data points, and the colored dotted lines are least-squares fits to the individual spectral types shown in the legend. For visual clarity, we show the size of the Wood et al. (2005) error bars in the gray-shaded box. Right:  $\text{Ly}\alpha$  surface flux vs. stellar rotation period. Similar to the left panel, the solid points and lines correspond to data from the MUSCLES sample, and the open symbols and dotted lines correspond to data from Wood et al. (2005).

(slope,  $p$ -value,  $\sigma$ ) are similar to those for Si III and C IV. This might seem surprising for C II, because of its similar formation temperature to  $\text{Ly}\alpha$ , but Oranje (1986) noted that C II behaves like a transition region line in M dwarfs. Fontenla et al. (2016) provide support of C II as a transition region line with their semi-empirical model of GJ 832’s atmosphere, which shows that the transition region begins at temperatures  $\sim 2000$  K cooler than in the quiet Sun.

We find no significant correlation between the chromospheric line  $\text{Ly}\alpha$  and the transition region emission lines Si III, Si IV, C II, C IV, and He II, but a larger sample size is needed to place a strong constraint on a non-correlation. Oranje (1986), Schrijver & Rutten (1987), and Rutten et al. (1989) showed that M dwarfs show a deficiency in chromospheric emission compared to FGK stars, but equivalent amounts of transition region and coronal emission compared to FGK stars. This deficiency appears independent of stellar magnetic activity, but may increase toward later-type stars with cooler effective temperature and higher surface gravity (Rutten et al. 1989). A shallow-sloped correlation (compared to the slope of a correlation between two chromospheric lines) between a chromospheric and transition region line has only been shown for optically active M dwarfs (e.g., Oranje 1986; albeit with sample sizes similar to the MUSCLES M-dwarf sample).

We use the  $\text{Ly}\alpha$ – $\text{Mg II}$  best-fit line (Figure 8) to correct GJ 1214’s  $\text{Ly}\alpha$  flux. As discussed in Section 4.7, GJ 1214 is a weak  $\text{Ly}\alpha$  emitter and its low-S/N spectrum may suffer from oversubtraction of airglow. By eye, it appears that our solution underestimates the  $\text{Ly}\alpha$  flux, and we are not confident in the robustness of its  $\text{Ly}\alpha$  profile reconstruction. Using the linear fit to the  $\text{Ly}\alpha$ – $\text{Mg II}$  correlation (which excludes GJ 1214), we increase the total flux of GJ 1214 by a factor of 2.4. Figure 8 shows the MCMC best-fit reconstructed flux for GJ 1214 labeled as “original” and the scaled flux labeled as “adjusted.”

In Figure 9 we compare our observed  $\text{Ly}\alpha$ – $\text{Mg II}$  and  $\text{Ly}\alpha$ –rotation period correlations to those that Wood et al. (2005) found for stars of spectral classes F V–M V. The MUSCLES targets probe a new parameter space of lower-mass and less-active stars. The addition of the MUSCLES K dwarfs to Wood et al.’s (2005) K-dwarf sample skews the relation to expect greater  $\text{Ly}\alpha$  fluxes for given  $\text{Mg II}$  flux on the low-activity end of the parameter space. However, this difference in slope is not statistically significant, and the low-activity part of the parameter space has not been sampled with K-dwarf observations. Wood et al.’s (2005) sample of four M dwarfs had only two with both  $\text{Ly}\alpha$  and  $\text{Mg II}$  flux measurements (AD Leo and AU Mic), and they could not create a M dwarf  $\text{Ly}\alpha$ – $\text{Mg II}$  relation. The addition of these two stars to the MUSCLES sample nearly doubles the slope of the  $\text{Ly}\alpha$ – $\text{Mg II}$  relation from what it would have been otherwise, but this difference is not statistically significant.

The addition of the MUSCLES K dwarfs to Wood et al.’s (2005) K dwarfs does not significantly change the  $\text{Ly}\alpha$ –rotation period relation. The entire Wood et al. (2005) M-dwarf sample (four stars) has rotation period measurements, and the addition of the MUSCLES M dwarfs (six stars excluding GJ 1214) allows the calculation of a  $\text{Ly}\alpha$ –rotation period correlation. The difference in slope between K and M dwarfs of the  $\text{Ly}\alpha$ –rotation period correlations is statistically significant at the  $2.4\sigma$  level. The correlations suggest that slowly rotating M dwarfs will exhibit less  $\text{Ly}\alpha$  flux than slowly rotating K dwarfs, but beyond the cross-over point of the relation, a faster-rotating M dwarf may exhibit more  $\text{Ly}\alpha$  flux than a similarly rotating K dwarf. However, this part of the parameter space has not been observationally explored for M dwarfs.

In Section 3.6, we utilized empirical relations from Linsky et al. (2014) to estimate the broadband extreme-UV spectra from the reconstructed, intrinsic  $\text{Ly}\alpha$  fluxes for the MUSCLES



sample. The Ly $\alpha$  and extreme-UV fluxes should be positively correlated, because they both depend on the magnetic heating rate. Models of solar quiet and active regions show similarly shaped temperature-pressure profiles (Fontenla et al. 2013), so all flux diagnostics should scale together, although not necessarily at the same rate. Claire et al. (2012) showed that the fraction of Ly $\alpha$  photons with respect to all photons at  $\lambda < 1700 \text{ \AA}$  is  $\sim 40\%$  throughout the Sun's history, and this supports Ly $\alpha$  flux as a good predictor of the extreme-UV flux. Much of the 100–1170  $\text{\AA}$  wavelength range contains H I and He I continua and emission lines with similar formation temperatures to Ly $\alpha$ . However, Ly $\alpha$  is probably not a good predictor of the higher formation temperature transition region and coronal emission lines that dominate the 100–400  $\text{\AA}$  region of the extreme-UV, as seen by the poor correlations with transition region ions such as C IV. However, extreme-UV reconstructions based on coronal X-rays and reconstructions based on chromospheric Ly $\alpha$  agree well, within  $\sim 30\%$  (Section 4.3).

For the Sun, de Wit et al. (2005) did not find that Ly $\alpha$  is the best emission line to predict the extreme-UV spectrum. If using only one or two emission lines to predict the solar extreme-UV spectrum, they recommend using O I ( $\lambda 1304 \text{ \AA}$ ), C IV ( $\lambda 1550 \text{ \AA}$ ), or Si II ( $\lambda 1817 \text{ \AA}$ ), or Ly $\alpha$  in conjunction with Si II ( $\lambda 1817 \text{ \AA}$ ) or O I ( $\lambda 1304 \text{ \AA}$ ). Models of the chromospheres, transition regions, and coronae of M and K dwarfs, using the line fluxes provided by the MUSCLES survey, will be required to identify more robust scaling relations for low-mass stars.

## 6. SUMMARY

Using an MCMC technique, we have reconstructed the intrinsic Ly $\alpha$  profiles from the observed, attenuated profiles for 11 low-mass exoplanet host stars, and used the intrinsic Ly $\alpha$  fluxes to estimate the extreme-UV spectra. We assume an intrinsic Ly $\alpha$  profile consisting of a dominant narrow Gaussian and a weak broad Gaussian, and a single interstellar absorption component characterized by a Voigt profile. Our results are summarized as follows.

1. The intrinsic Ly $\alpha$  profiles are broad ( $\sim 300 \text{ km s}^{-1}$ ) with wings that extend to about  $\pm 1200 \text{ km s}^{-1}$  ( $\Delta\lambda \approx 5 \text{ \AA}$ ). The vast majority of the flux is in the profile's narrow component; the wings beyond  $\pm 300 \text{ km s}^{-1}$  contain at least 1% of the total flux.
2. The intrinsic Ly $\alpha$  surface flux correlates positively with Mg II, another chromospheric emission line, for both K and M dwarfs. There may be positive correlations of the Ly $\alpha$  surface flux with lines formed at higher temperatures in the stellar atmosphere (e.g., Si III, C IV), but our K and M dwarf samples are too small to make a statistically significant statement at this time.
3. We have found new Ly $\alpha$ –Mg II surface flux and Ly $\alpha$ –rotation period relations for K and M spectral types that span active and inactive stars. The relations change between spectral type probably owing to different temperature-pressure profiles in the stellar atmospheres.
4. We measure H I column densities for 10 new sightlines ranging from 4.7–21.1 pc distant. The lower spectral resolution of the G140M data does not influence H I column densities inferred from the higher-resolution E140M spectra. Log $_{10}$  H I column densities range from 18.0–18.6 with one sightline, GJ 176, having an

extremely low log $_{10} N(\text{H I}) = 17.5$ . GJ 176 would make an excellent target for future extreme-UV observations.

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This research made use of the Python packages *emcee* (Foreman-Mackey et al. 2013), *pyspeckit* (Ginsburg & Mirocha 2011), *lyapy* (<https://github.com/allison/lyapy>), *scicatalog* (<https://github.com/parkus/scicatalog>), and *triangle* (Foreman-Mackey et al. 2014).

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